

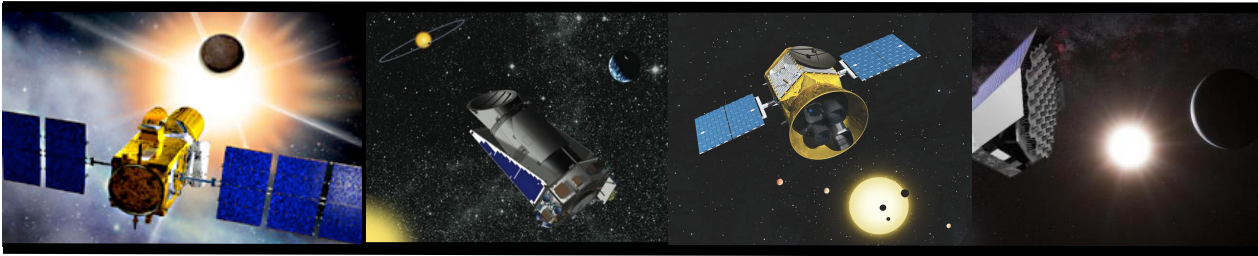
Tidal dissipation in stars and planetary fluid layers

S. Mathis, P. Auclair-Desrotour,
M. Guenel, C. Le Poncin-Lafitte



General context

The space photometry revolution: the discovery of new planetary systems and the characterisation of their host stars



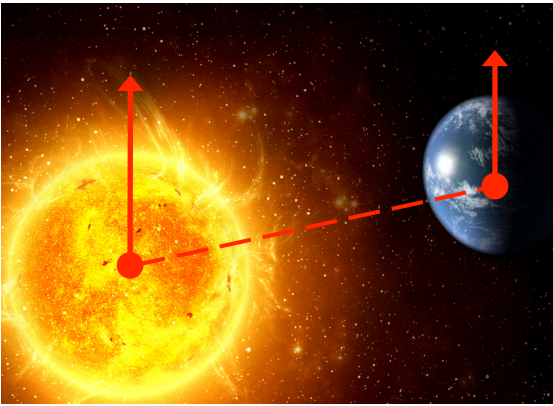
CoRoT

Kepler – K2

CHEOPS & TESS

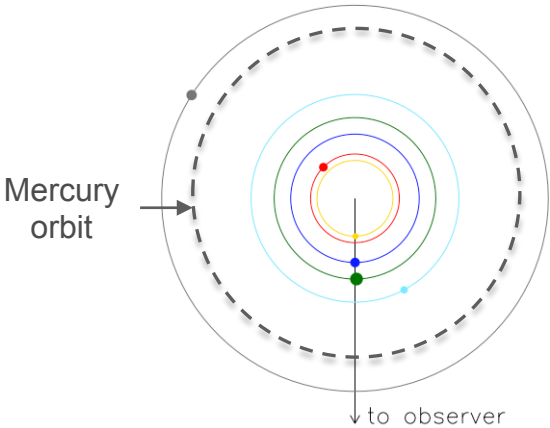
PLATO

Stellar and planetary rotation history



Albrecht et al. 2012; Gizon et al. 2013;
Talks C. Damiani & T. Mazey

Orbital architecture

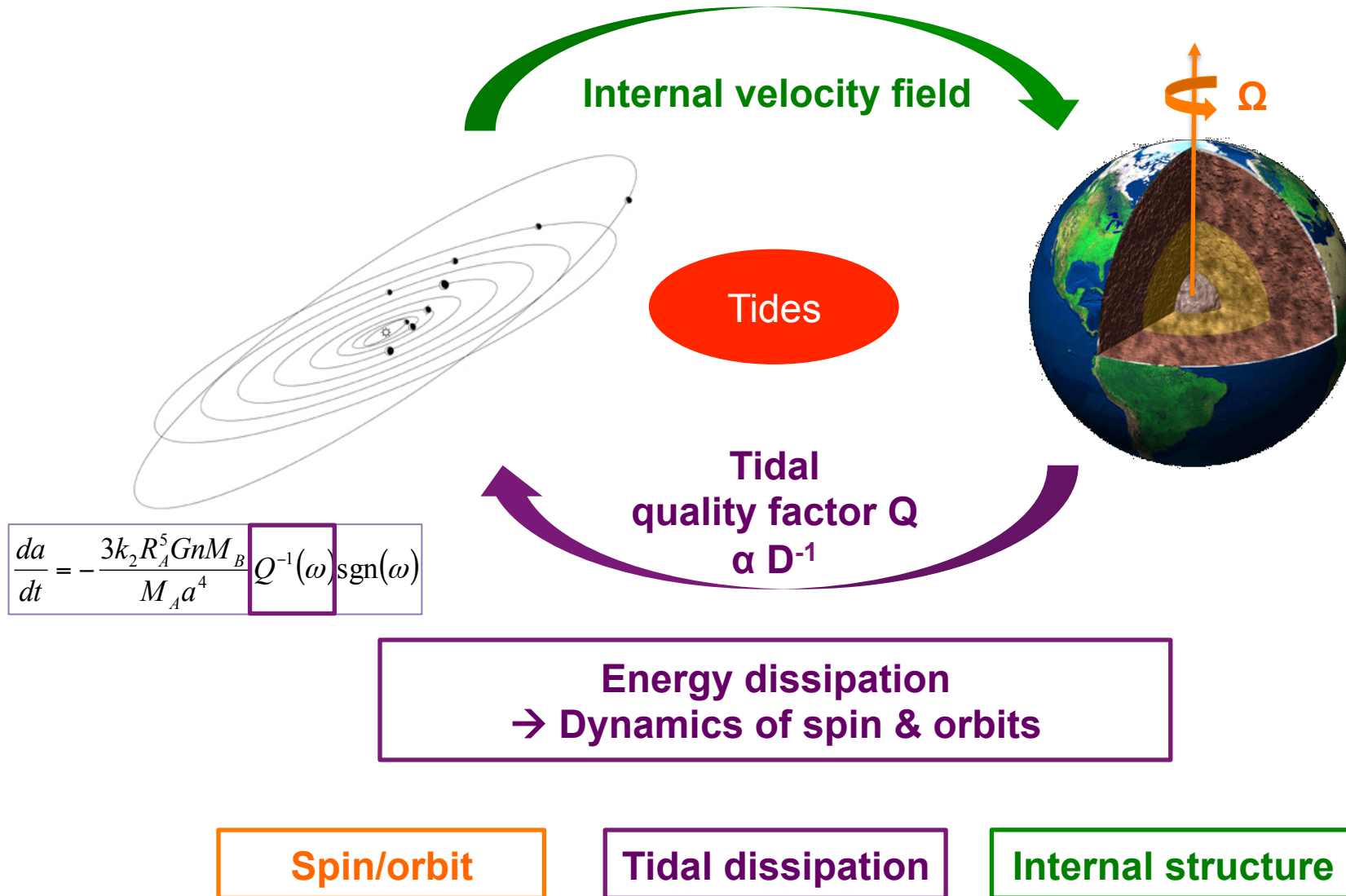


Kepler 11

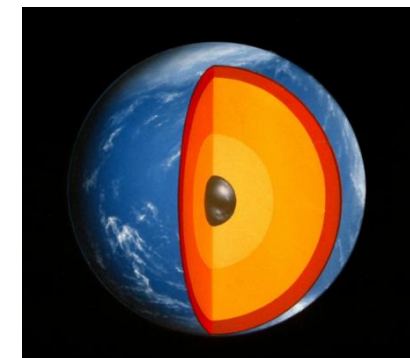
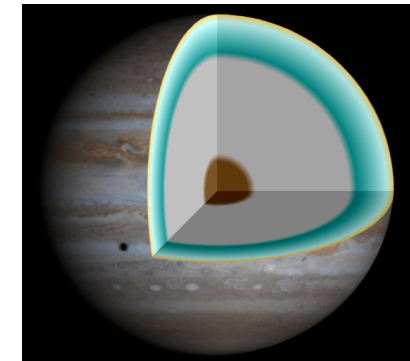
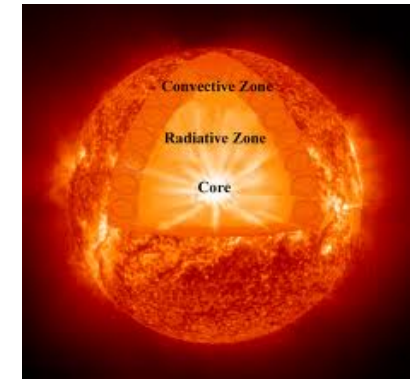
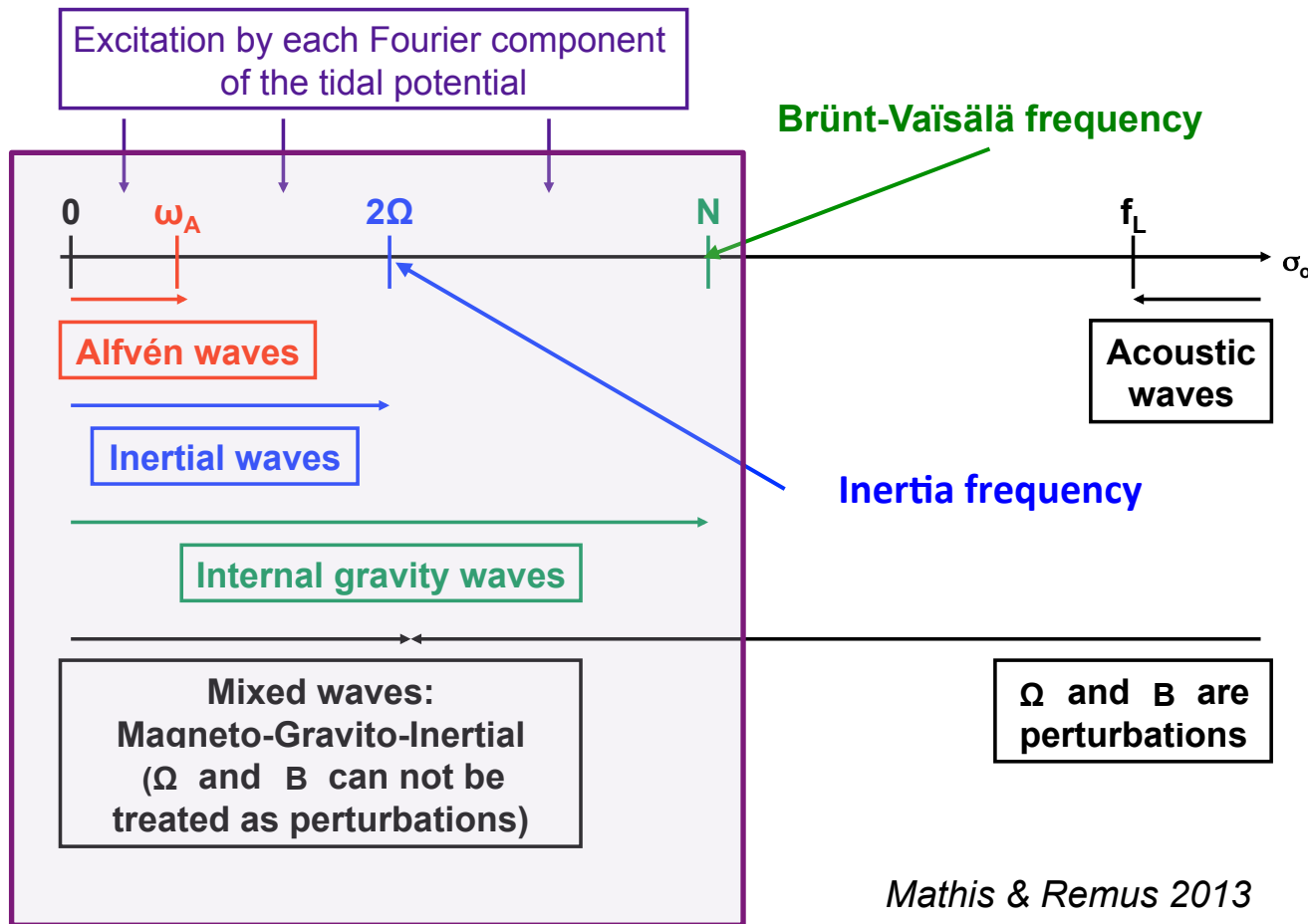
Lissauer et al. 2011

→ Need to understand angular momentum exchanges within star-planet systems → TIDES

Tidal dissipation

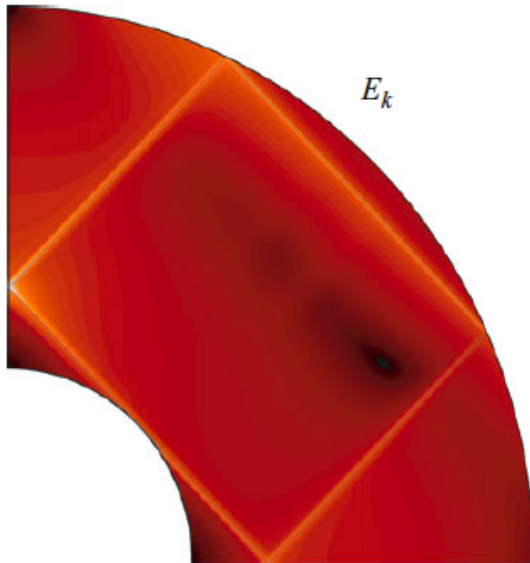
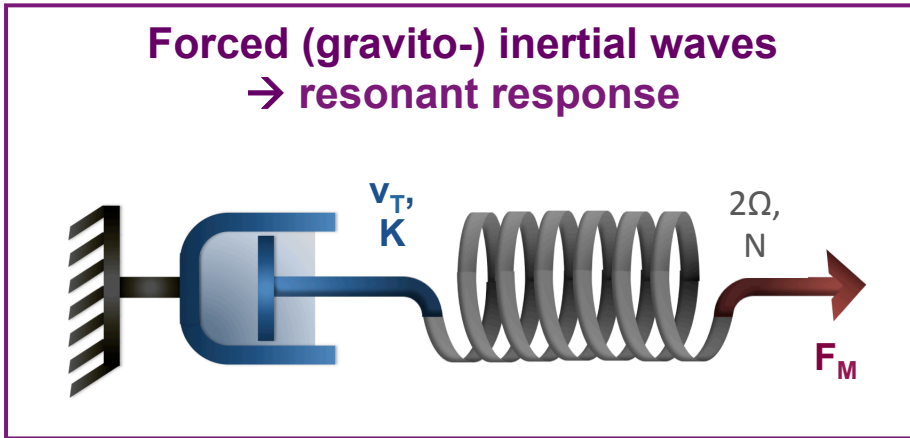


Tidal waves in stars and fluid planetary layers

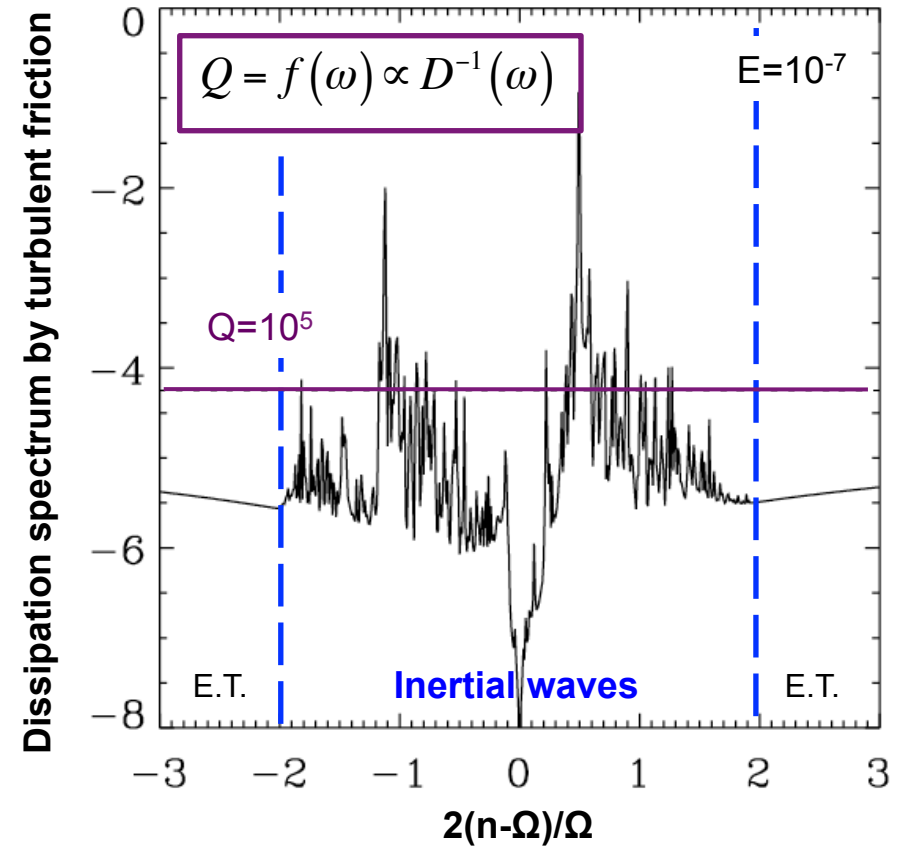


- **Convective layers:** turbulent friction & thermal diffusion
- **Stable layers:** thermal diffusion

A resonant erratic tidal dissipation spectrum



$Nr = 400 \quad L = 1200 \quad M = 0^+ \quad E = 2.0 \times 10^{-9}$
 $\eta = 0.350 \quad \omega = 0.707 \quad CL = ff$

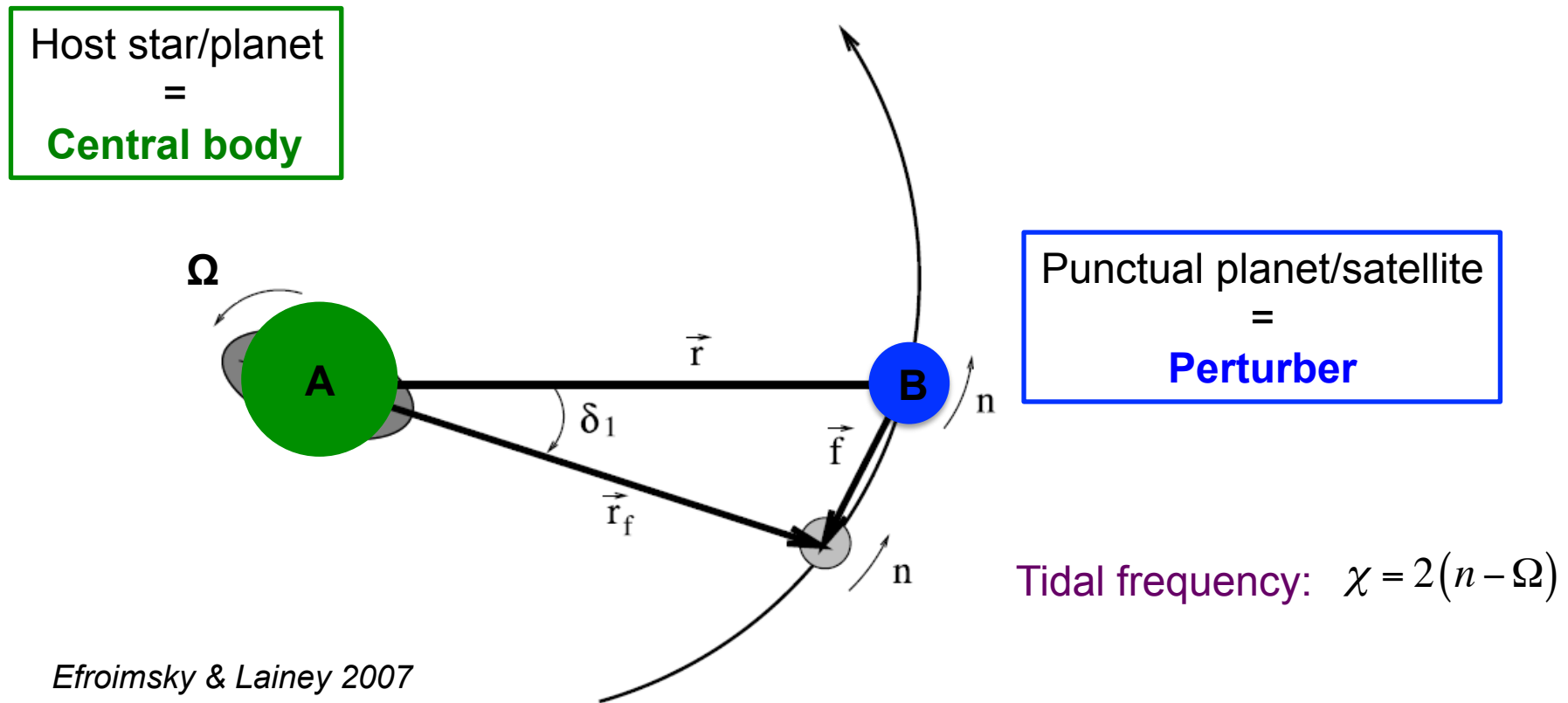


Ogilvie & Lin 2004: the case of Jupiter

Rieutord & Valdettaro 2010
 See also Ogilvie & Lin 2007

The impact of tidal dissipation on the spin dynamics and on systems orbital architecture

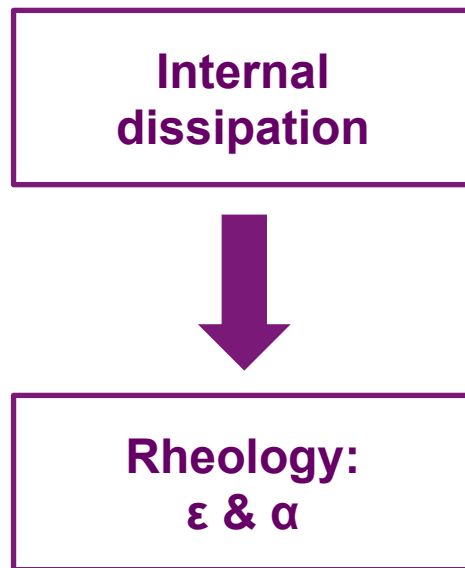
The coplanar two-bodies system:



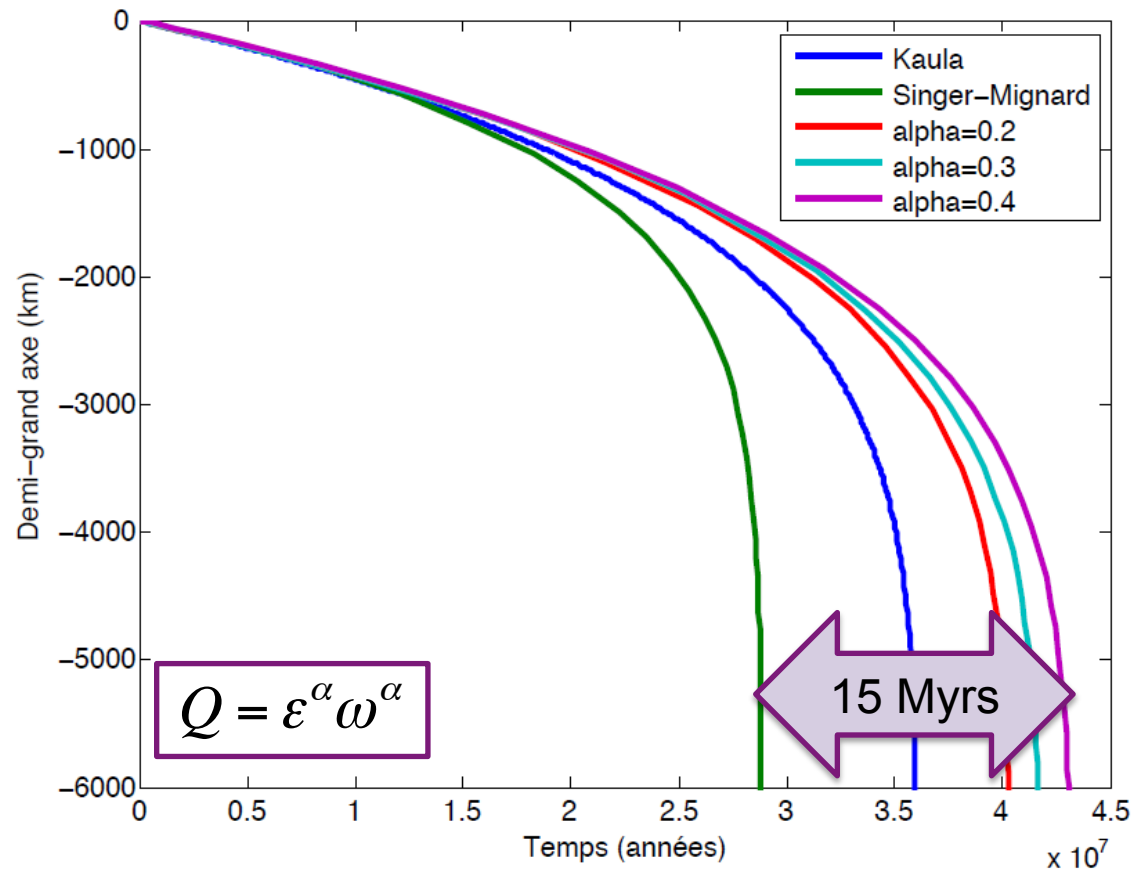
The impact of tidal dissipation: the case of rocky bodies

An example: the Mars-Phobos system

Regular evolution



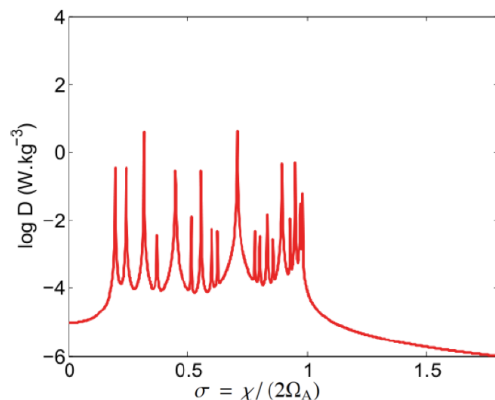
Efroimsky & Lainey 2007



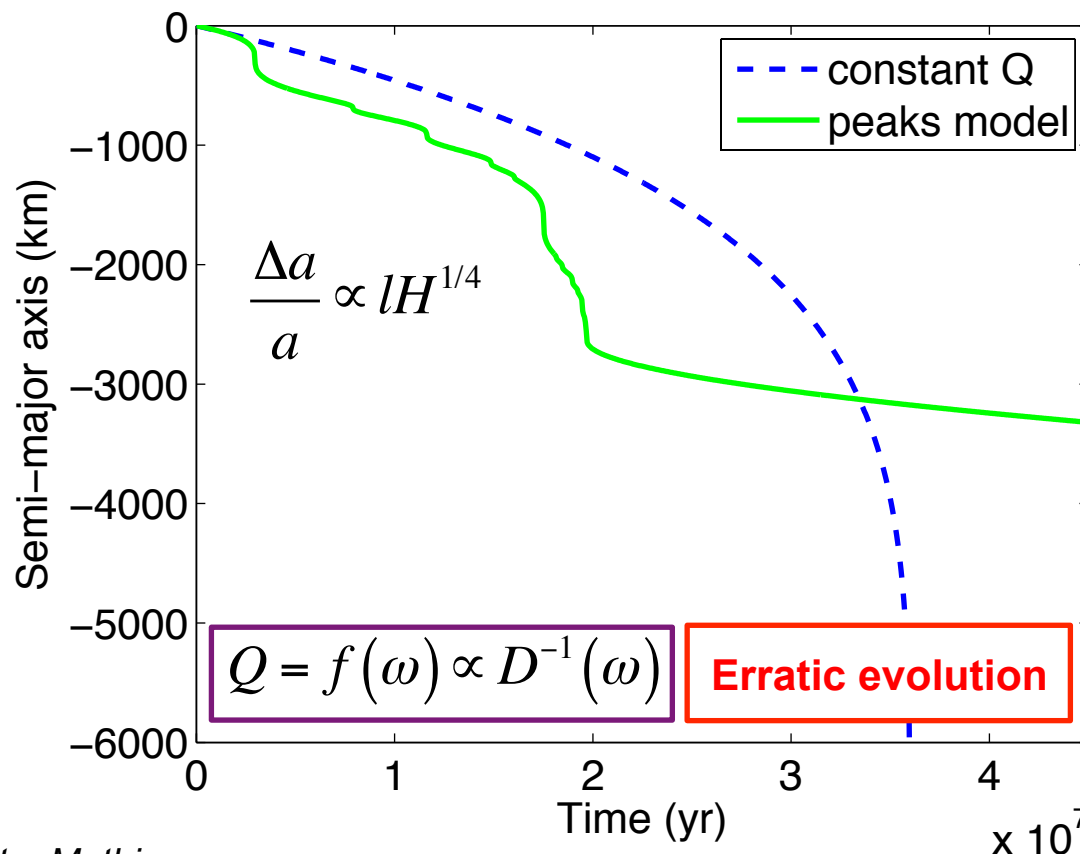
The impact of tidal dissipation: the case of fluid bodies

An example: a fully convective body of the mass of Mars-Phobos system

**Internal
dissipation**

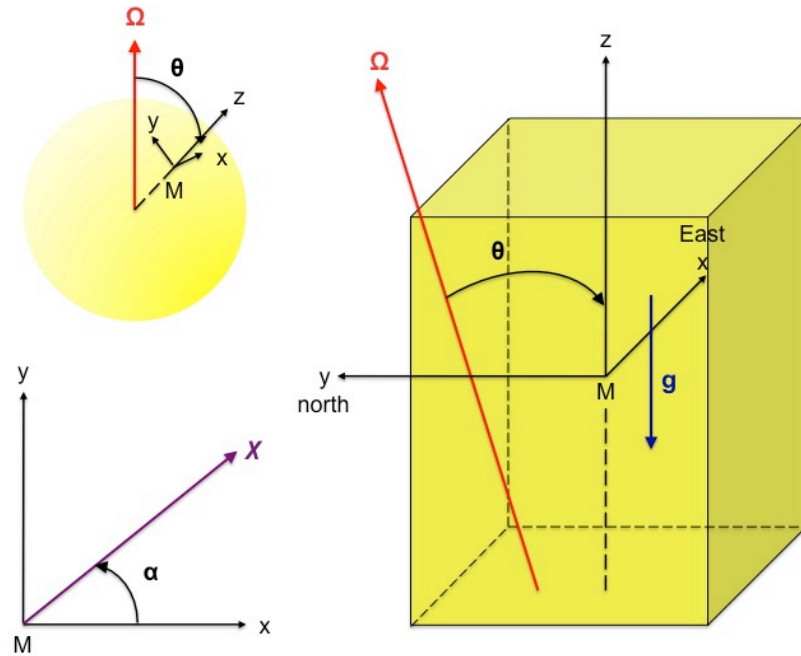


Fluid parameters



Auclair-Desrotour, Le Poncin-Lafitte, Mathis
2014a

A reduced local model to understand tidal dissipation in fluids



Ogilvie & Lin 2004

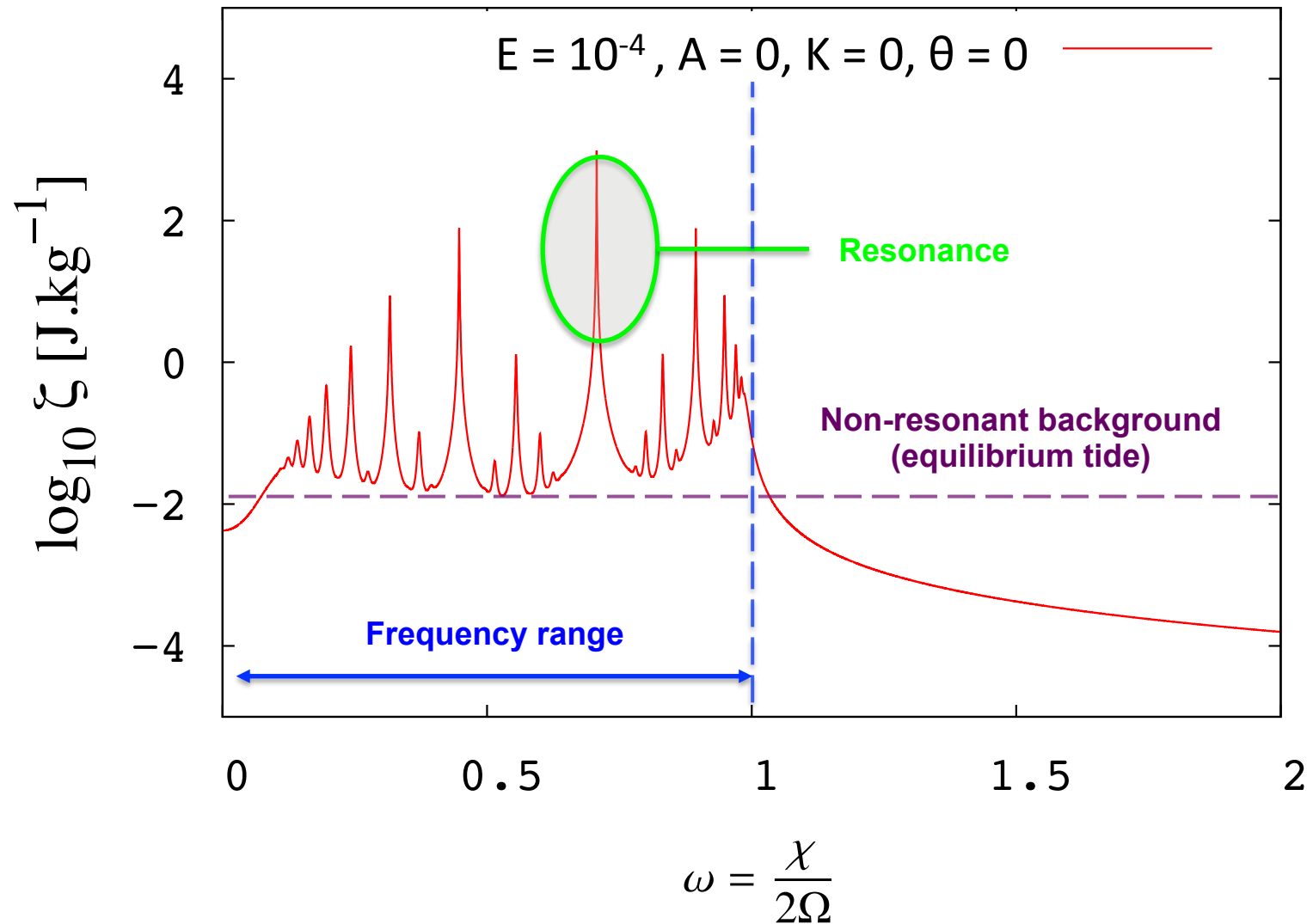
Auclair-Desrotour, Le Poncin-Lafitte, Mathis
2014b

- Cartesian geometry
- Rotating and **inclined**
- **Possible stable stratification**
- Viscous and **thermal dissipation**

Control parameters:

$A = \left(\frac{N}{2\Omega} \right)^2$	Stratification Coriolis
$E = \frac{2\pi^2\nu}{\Omega L^2}$	Viscous force Coriolis
$K = \frac{2\pi^2\kappa}{\Omega L^2}$	Thermal diffusivity Coriolis

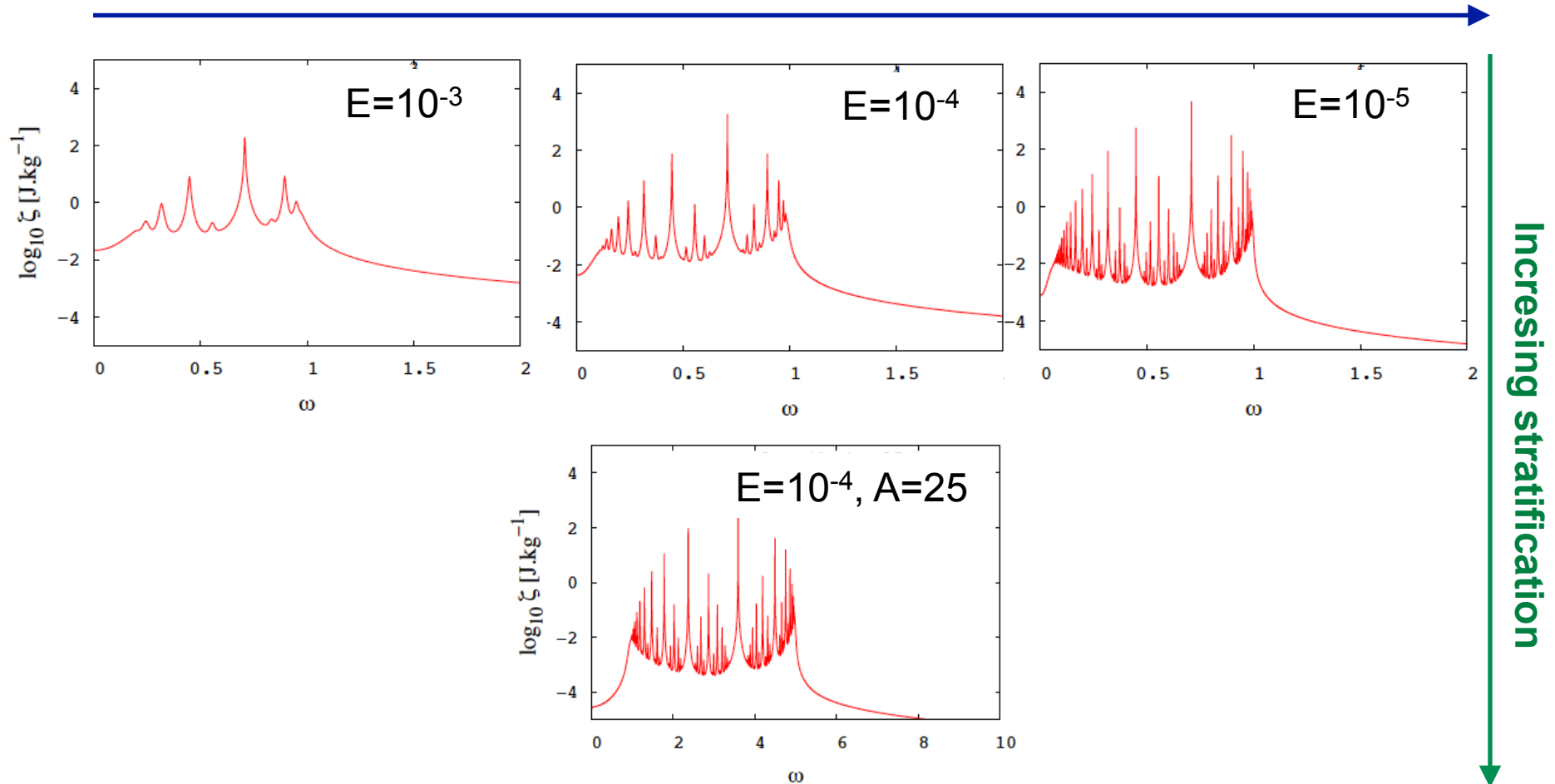
The complex erratic tidal dissipation spectrum



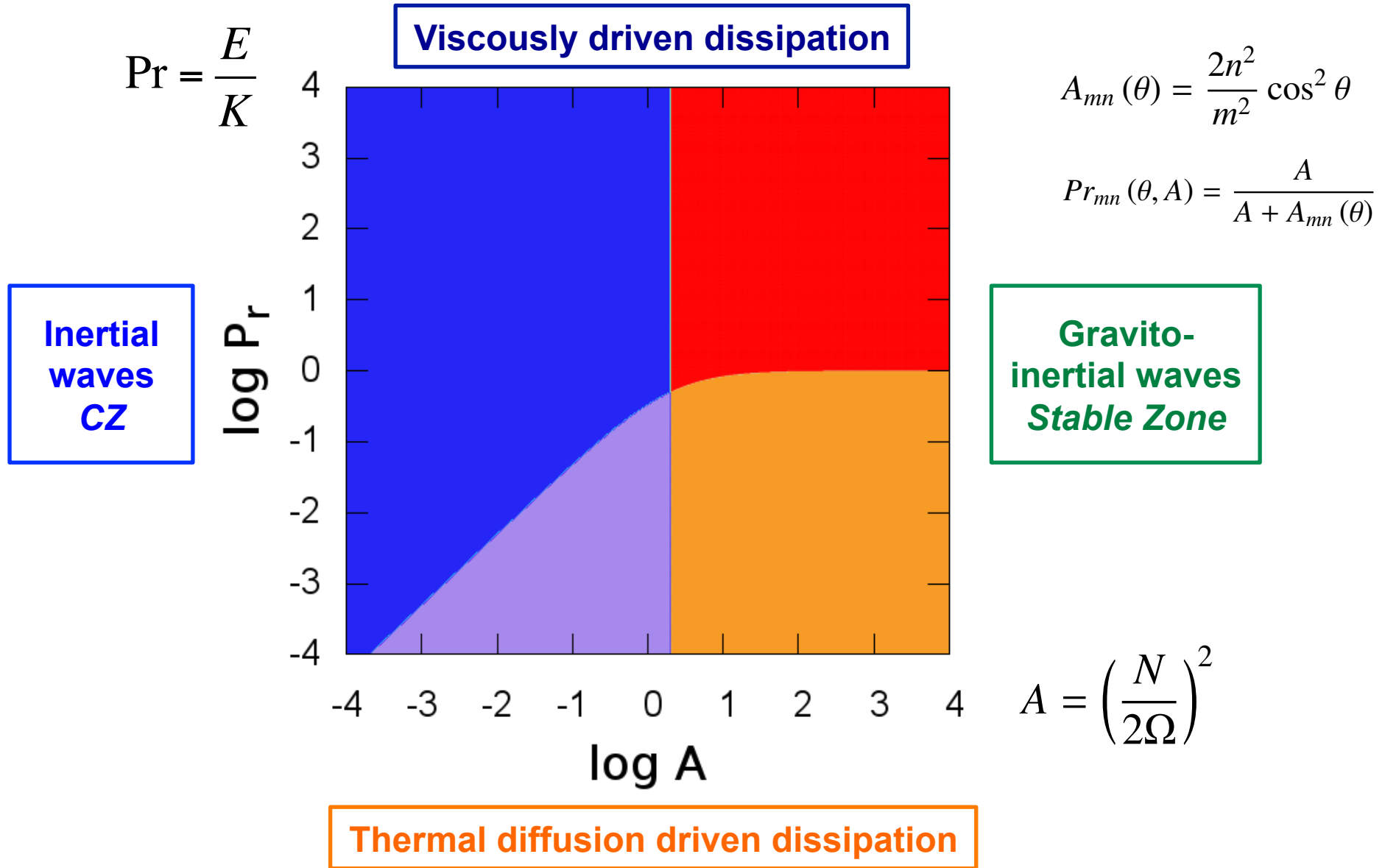
→ Need to characterize spectra

An evolving behaviour

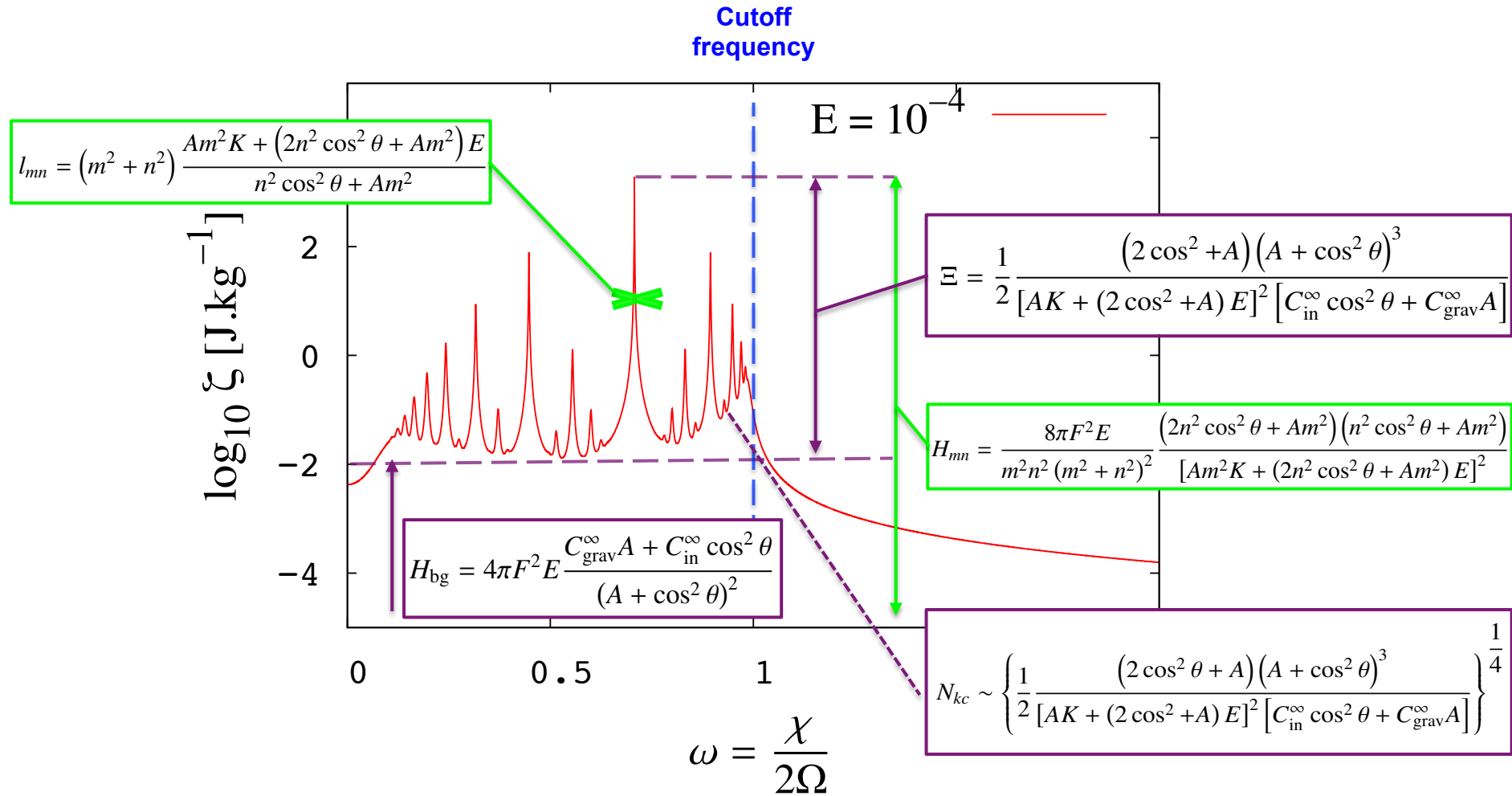
Decreasing viscosity / increasing rotation



The four main regimes



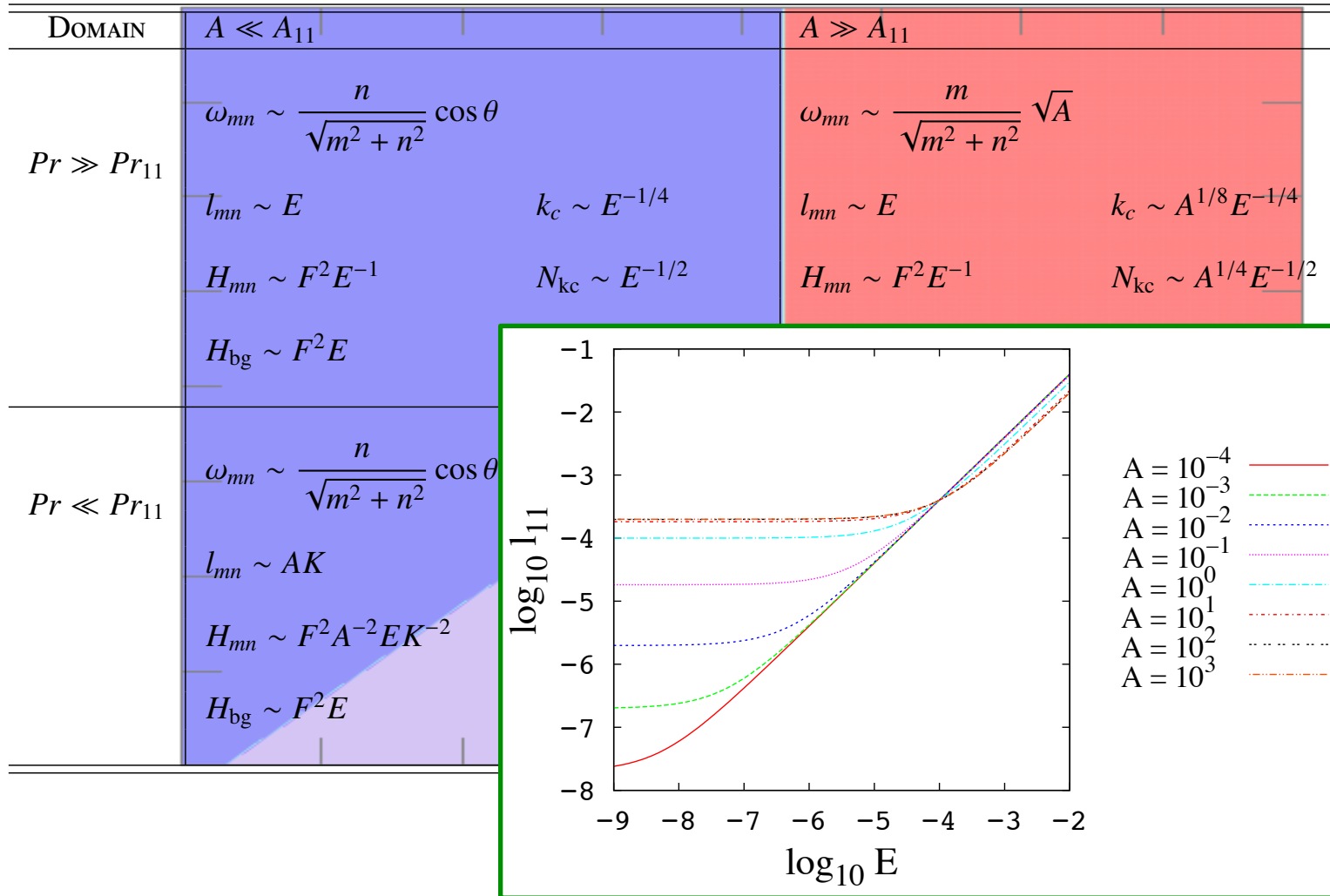
A complete characterization



Asymptotic scaling laws

DOMAIN	$A \ll A_{11}$	$A \gg A_{11}$
$Pr \gg Pr_{11}$	$\omega_{mn} \sim \frac{n}{\sqrt{m^2 + n^2}} \cos \theta$	$\omega_{mn} \sim \frac{m}{\sqrt{m^2 + n^2}} \sqrt{A}$
	$l_{mn} \sim E$ $k_c \sim E^{-1/4}$	$l_{mn} \sim E$ $k_c \sim A^{1/8} E^{-1/4}$
	$H_{mn} \sim F^2 E^{-1}$ $N_{kc} \sim E^{-1/2}$	$H_{mn} \sim F^2 E^{-1}$ $N_{kc} \sim A^{1/4} E^{-1/2}$
	$H_{bg} \sim F^2 E$ $\Xi \sim E^{-2}$	$H_{bg} \sim F^2 EA^{-1}$ $\Xi \sim AE^{-2}$
$Pr \ll Pr_{11}$	$\omega_{mn} \sim \frac{n}{\sqrt{m^2 + n^2}} \cos \theta$	$\omega_{mn} \sim \frac{m}{\sqrt{m^2 + n^2}} \sqrt{A}$
	$l_{mn} \sim AK$ $k_c \sim A^{-1/4} K^{-1/4}$	$l_{mn} \sim K$ $k_c \sim A^{1/8} K^{-1/4}$
	$H_{mn} \sim F^2 A^{-2} EK^{-2}$ $N_{kc} \sim A^{-1/2} K^{-1/2}$	$H_{mn} \sim F^2 EK^{-2}$ $N_{kc} \sim A^{1/4} K^{-1/2}$
	$H_{bg} \sim F^2 E$ $\Xi \sim A^{-2} K^{-2}$	$H_{bg} \sim F^2 EA^{-1}$ $\Xi \sim AK^{-2}$

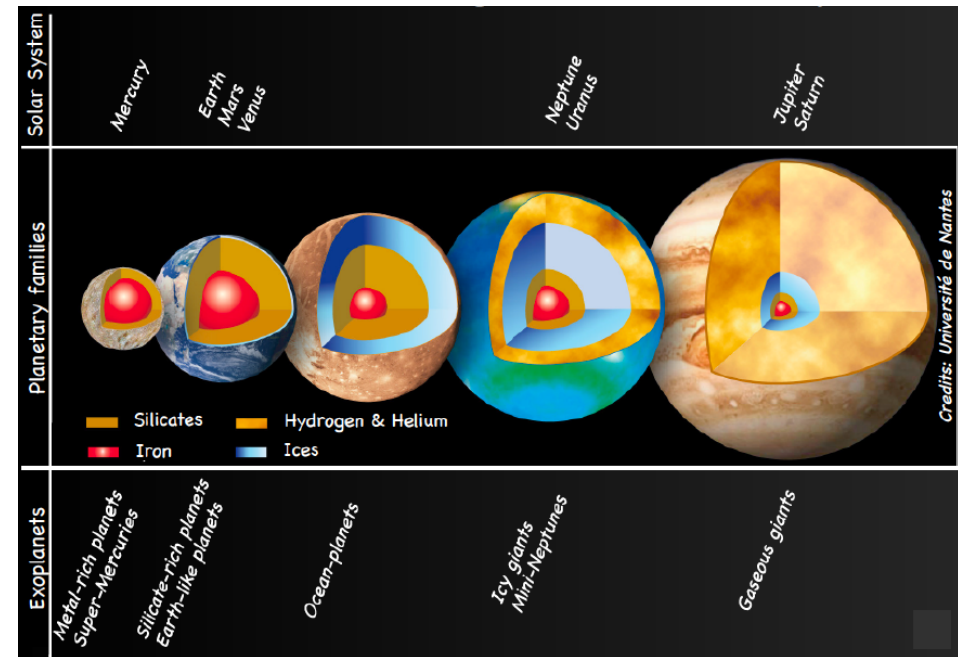
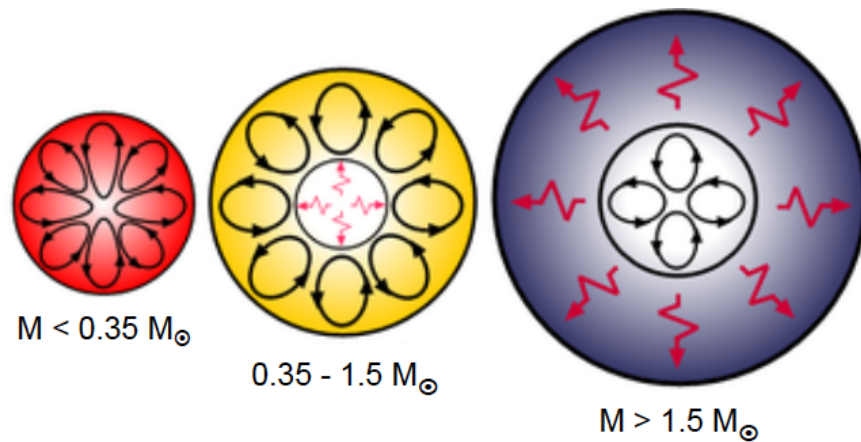
Asymptotic scaling laws



Towards global and multi-layer models

Host star (M in M_{\odot})

Planets



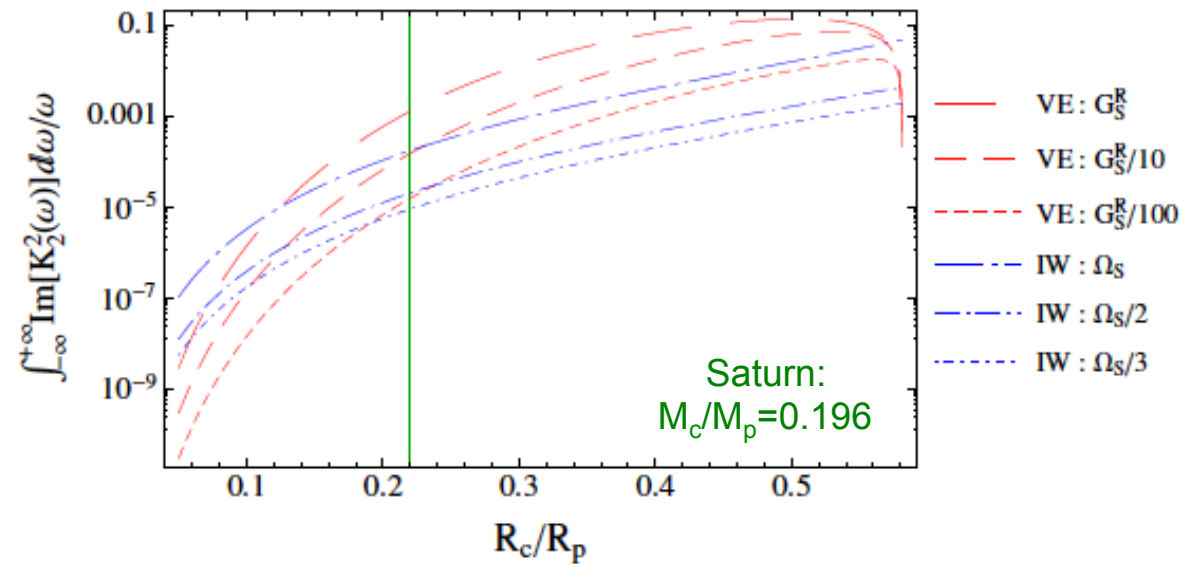
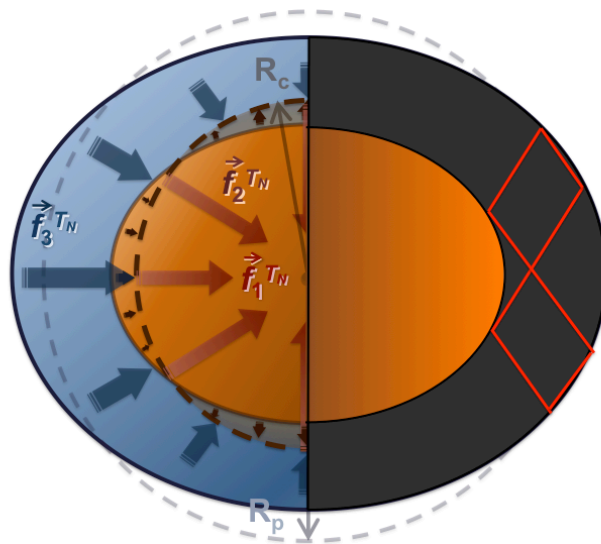
Their complex internal structure and rotation impact tidal dissipation
 → Need of an **ab-initio physical modeling**

Frequency-averaged models

The example of a Saturn-like planet:

Remus, Mathis,
Zahn & Lainey
2012

Ogilvie 2009, 2013



→ Integrated models needed for gaseous giant (and telluric) planets

→ Possibility of frequency-averaged grids as a function of stellar and planetary properties

Guenel, Mathis & Remus 2014

Conclusions & perspectives

- Dependence of the spin/orbital dynamics on the resonant tidal fluid dissipation :
→ *width, height, non-resonant background level*
- Dependence of the characteristics of these resonances on the physical parameters of the fluid :
→ *rotation, stratification, viscosity, thermal diffusivity, etc.*
- Local model : general method and qualitative results
→ *Need of global models (Guenel, Baruteau, Mathis & Rieutord; Ogilvie et al.); need to characterize the case of stratified convection (I. Baraffe's talk)*
- Generalization to magnetic stars and planets :
→ *Alfvén waves; new asymptotic behaviors (Mathis, Auclair-Desrotour, Guenel, Le Poncin-Lafitte)*

Spin/orbit

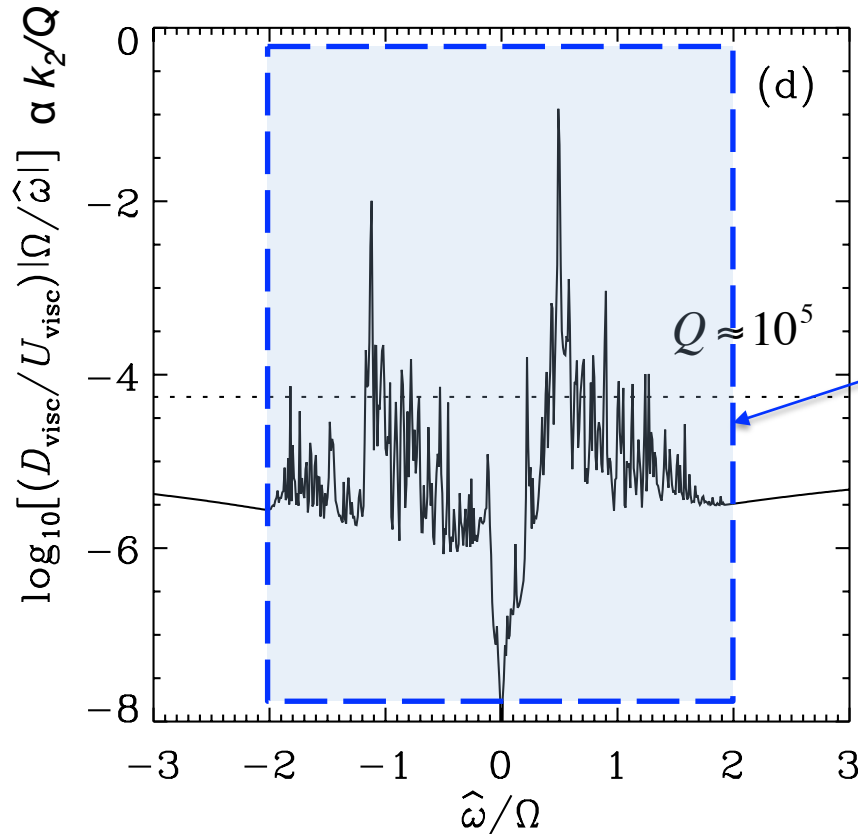
Tidal dissipation

Internal structure

APPENDIX

Tidal dissipation in stars and fluid planetary layers

→ A resonant erratic tidal dissipation spectrum



Tidal inertial waves
(Coriolis acceleration)

$$Q = f(\omega) \propto D^{-1}(\omega)$$

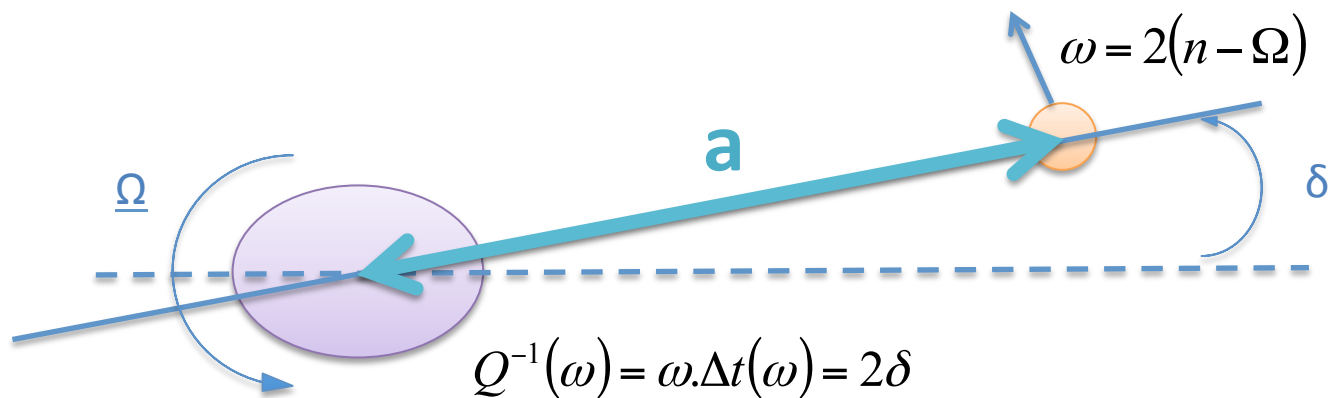
Ogilvie & Lin 2004: the case of Jupiter
See also Ogilvie & Lin 2007; Rieutord & Valdettaro 2010

II. What is the impact of tidal dissipation on the orbital evolution ?

- Orbital dynamics

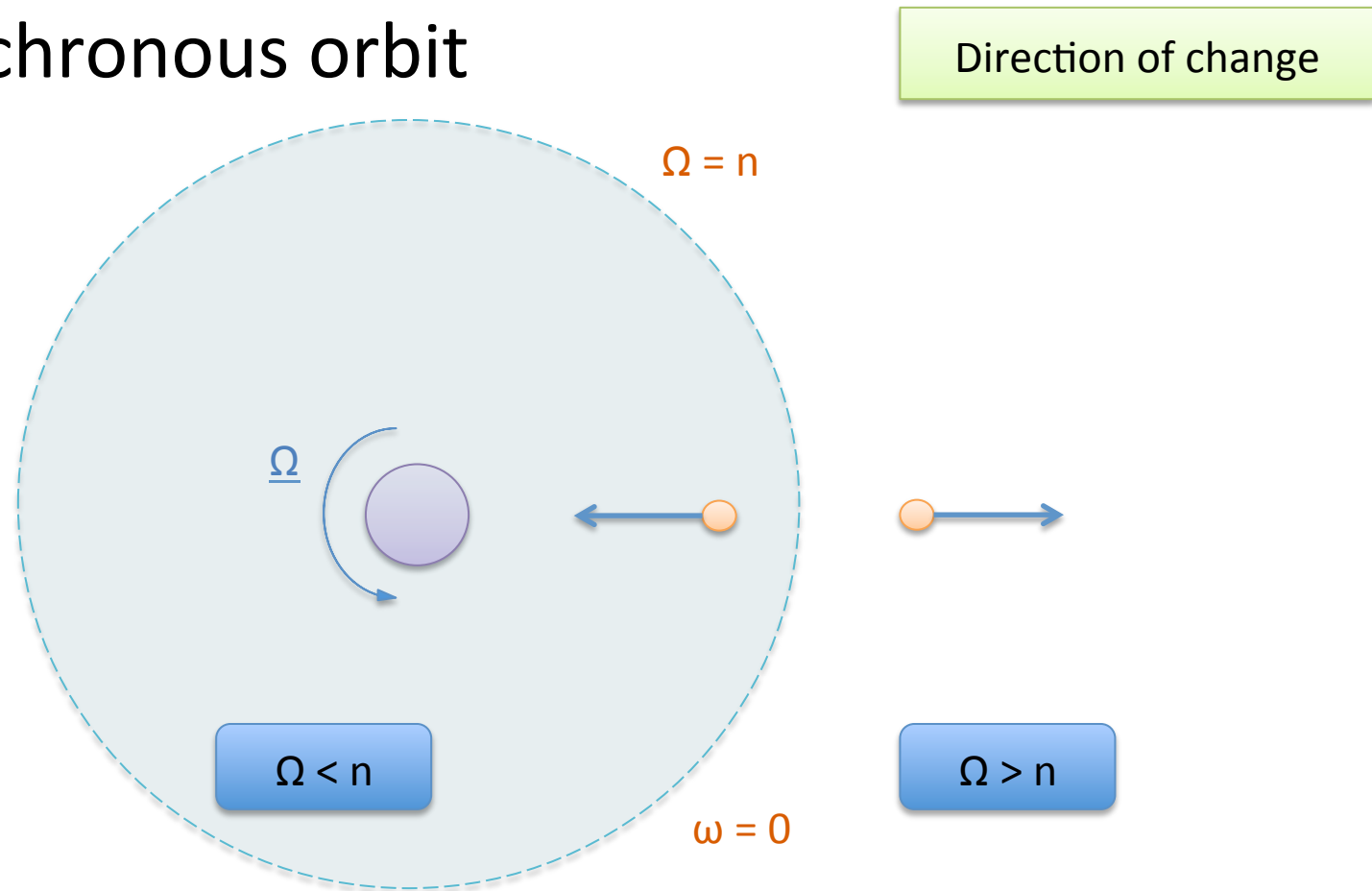
$$\frac{da}{dt} = -\frac{3k_2 R_A^5 G n M_B}{M_A a^4} Q^{-1}(\omega) \text{sgn}(\omega)$$

Relative motion
Direction of change
Velocity of change
Dissipation



II. What is the impact of tidal dissipation on the orbital evolution ?

- The synchronous orbit



II. What is the impact of tidal dissipation on the orbital evolution ?

- What about fluid bodies ?

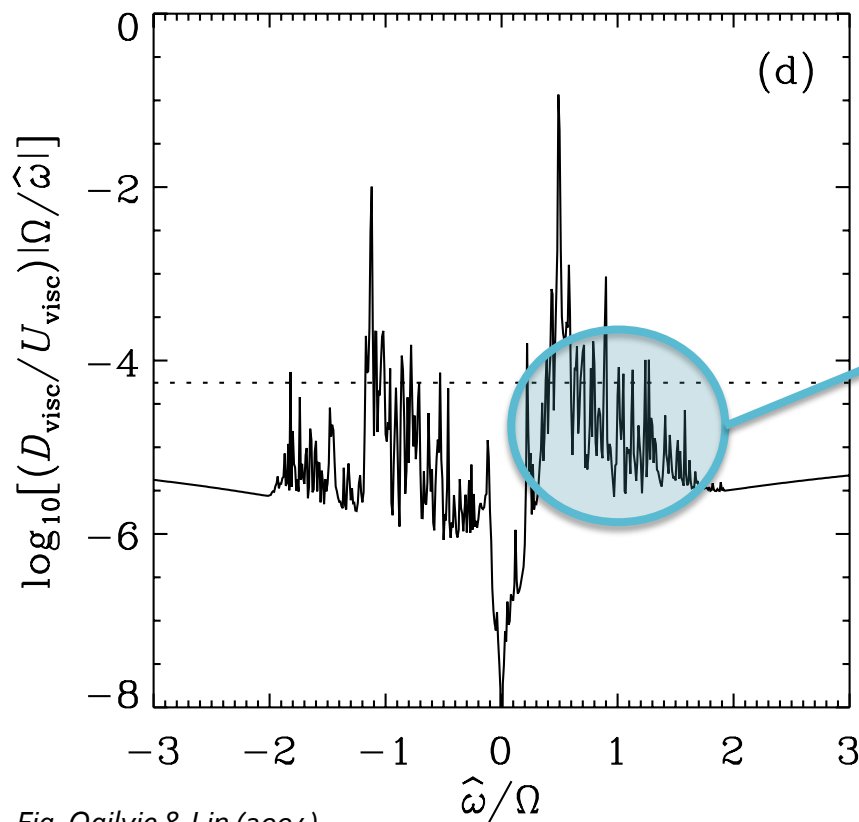


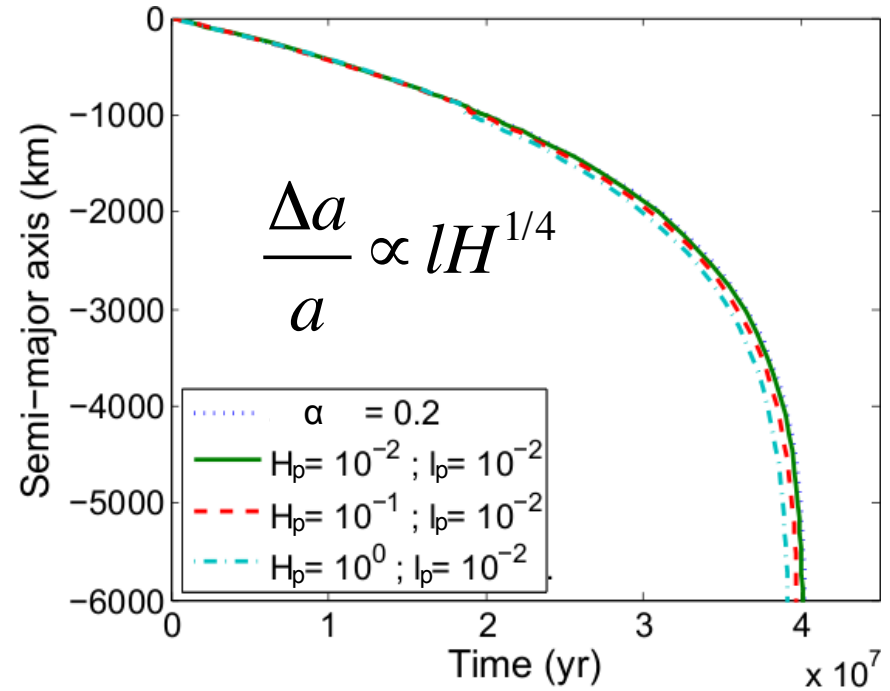
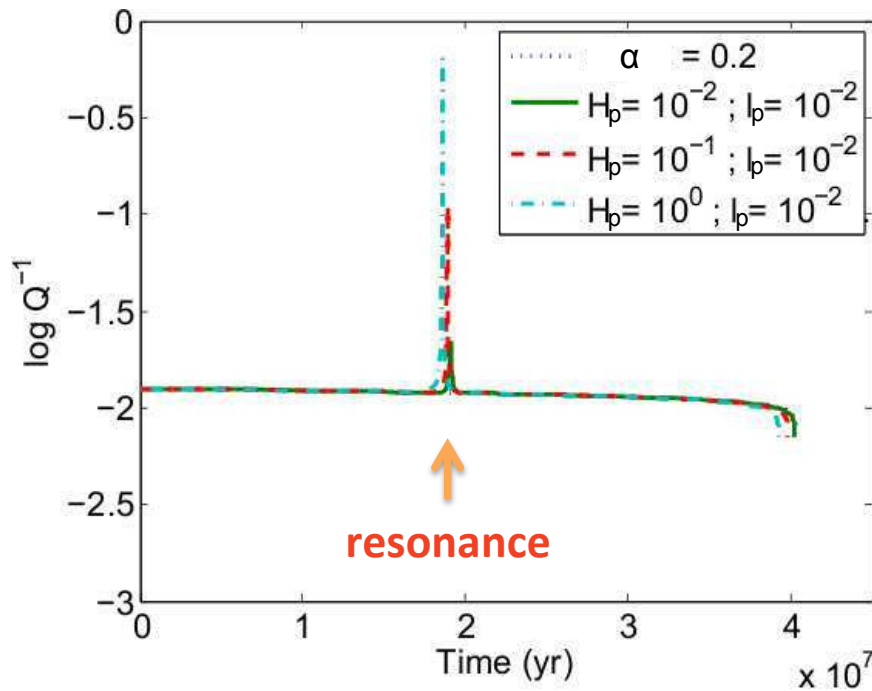
Fig. Ogilvie & Lin (2004)

Inertial waves
(Coriolis)

$$Q = f(\omega) \propto D^{-1}(\omega)$$

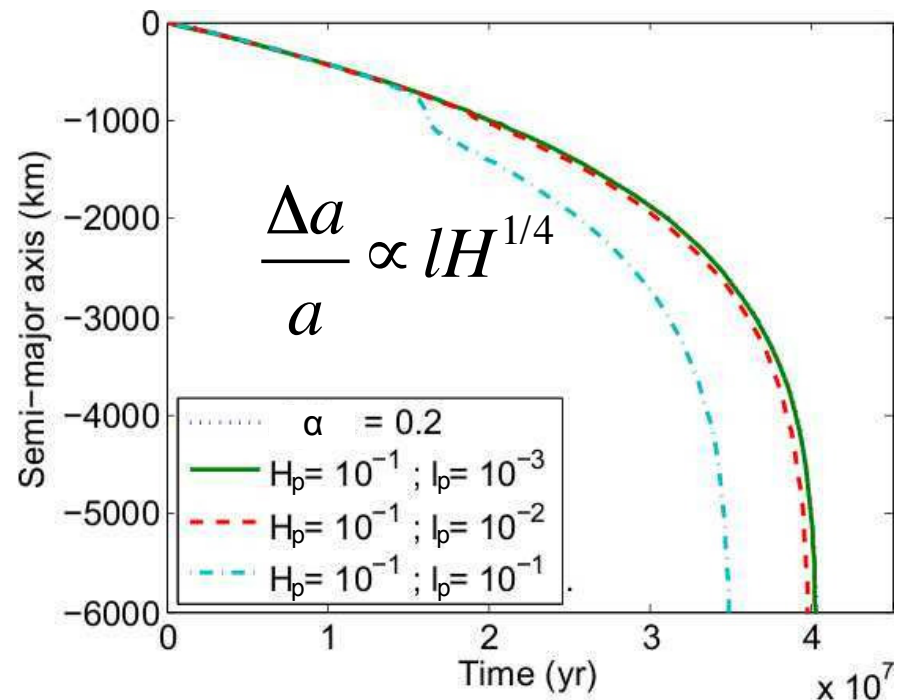
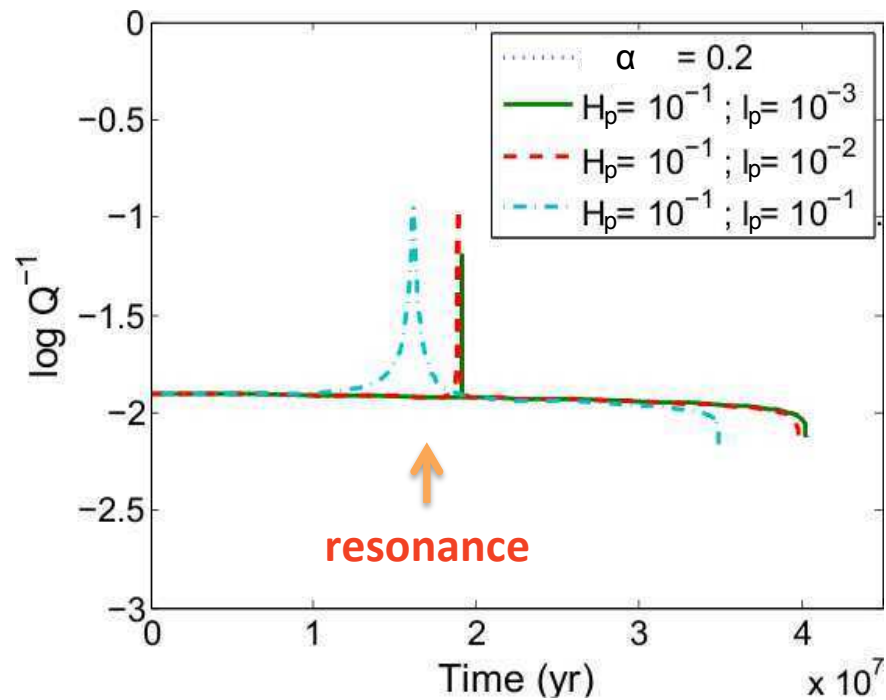
II. What is the impact of tidal dissipation on the orbital evolution ?

- Dependence on the height $Q_p^{-1}(\omega) = \frac{H_p}{\left[4(\sqrt{2}-1)\left(\frac{\omega-\omega_p}{l_p}\right)^2 + 1\right]^2}$



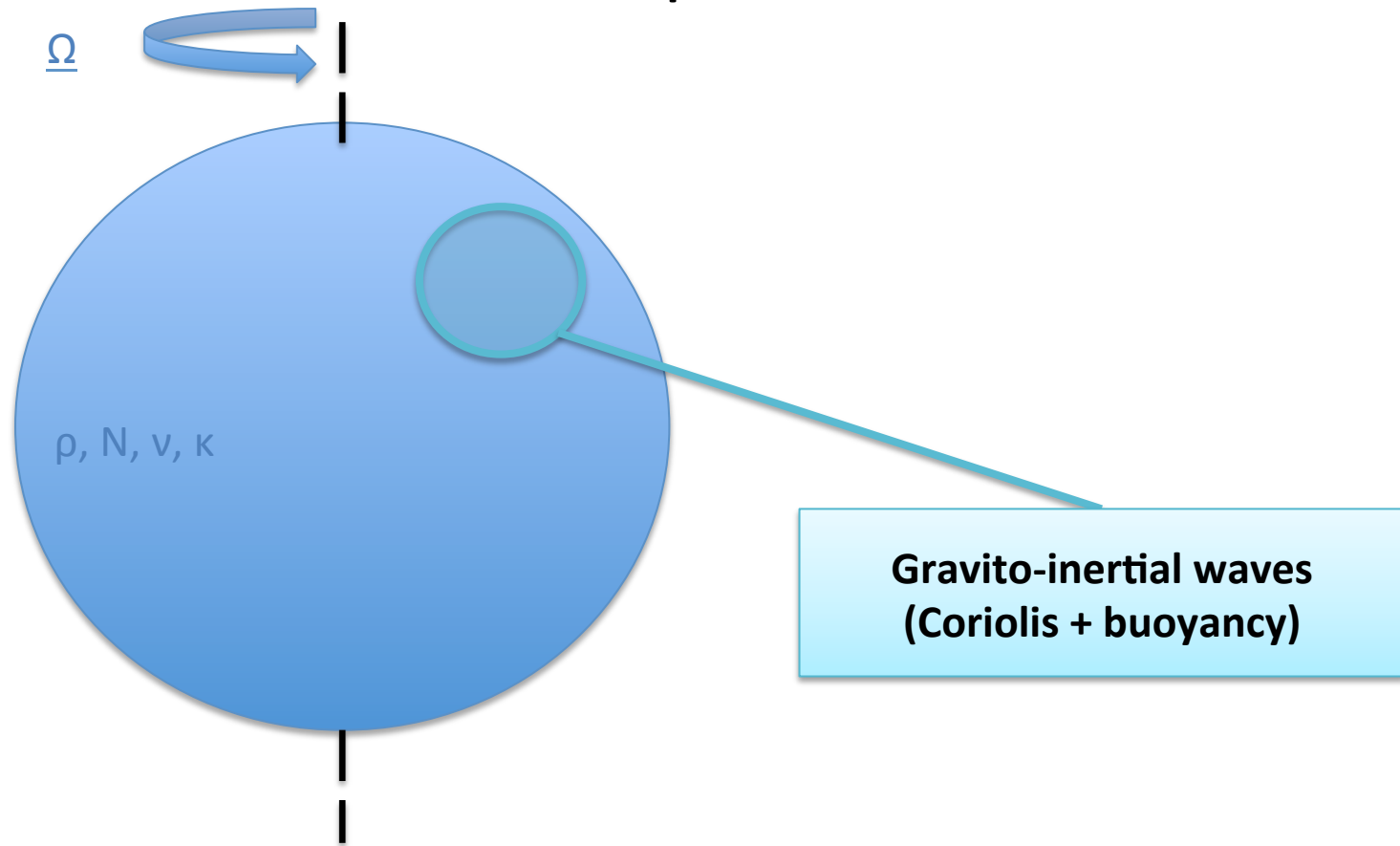
II. What is the impact of tidal dissipation on the orbital evolution ?

- Dependence on the width $Q_p^{-1}(\omega) = \frac{H_p}{\left[4(\sqrt{2}-1)\left(\frac{\omega-\omega_p}{l_p}\right)^2 + 1\right]^2}$



III. How does tidal dissipation depend on the internal physics of the bodies ?

- A local model for a fluid planet



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- A local model for a fluid planet

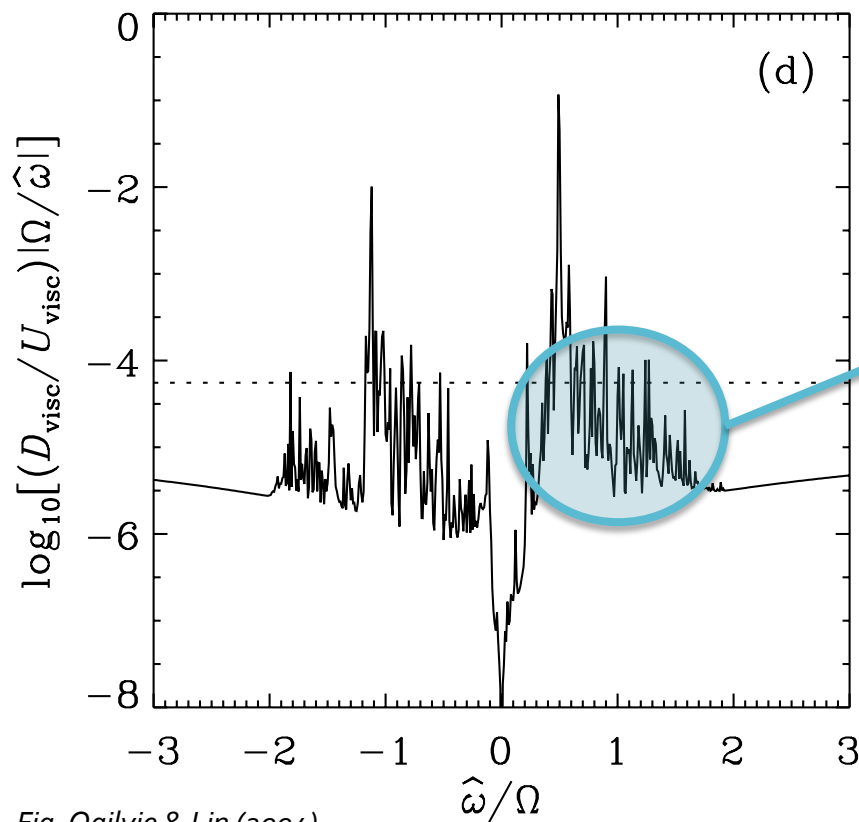
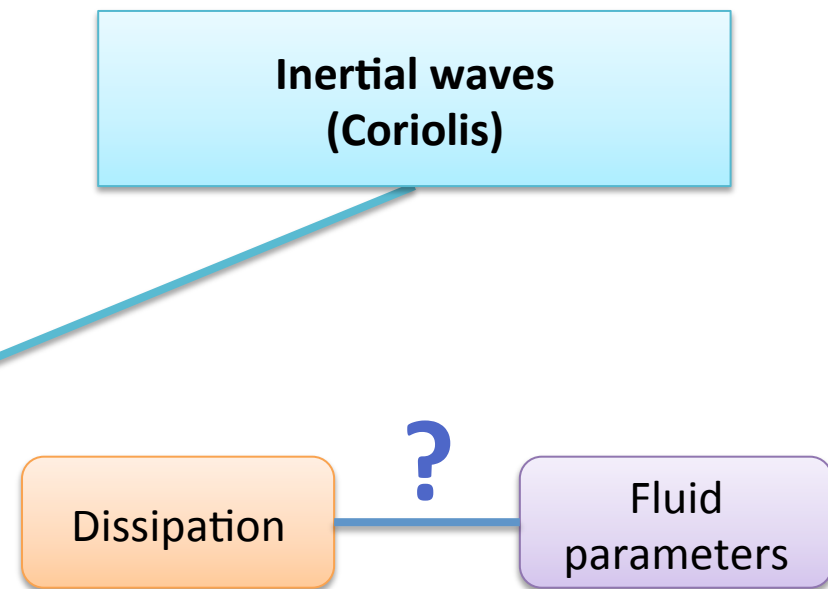
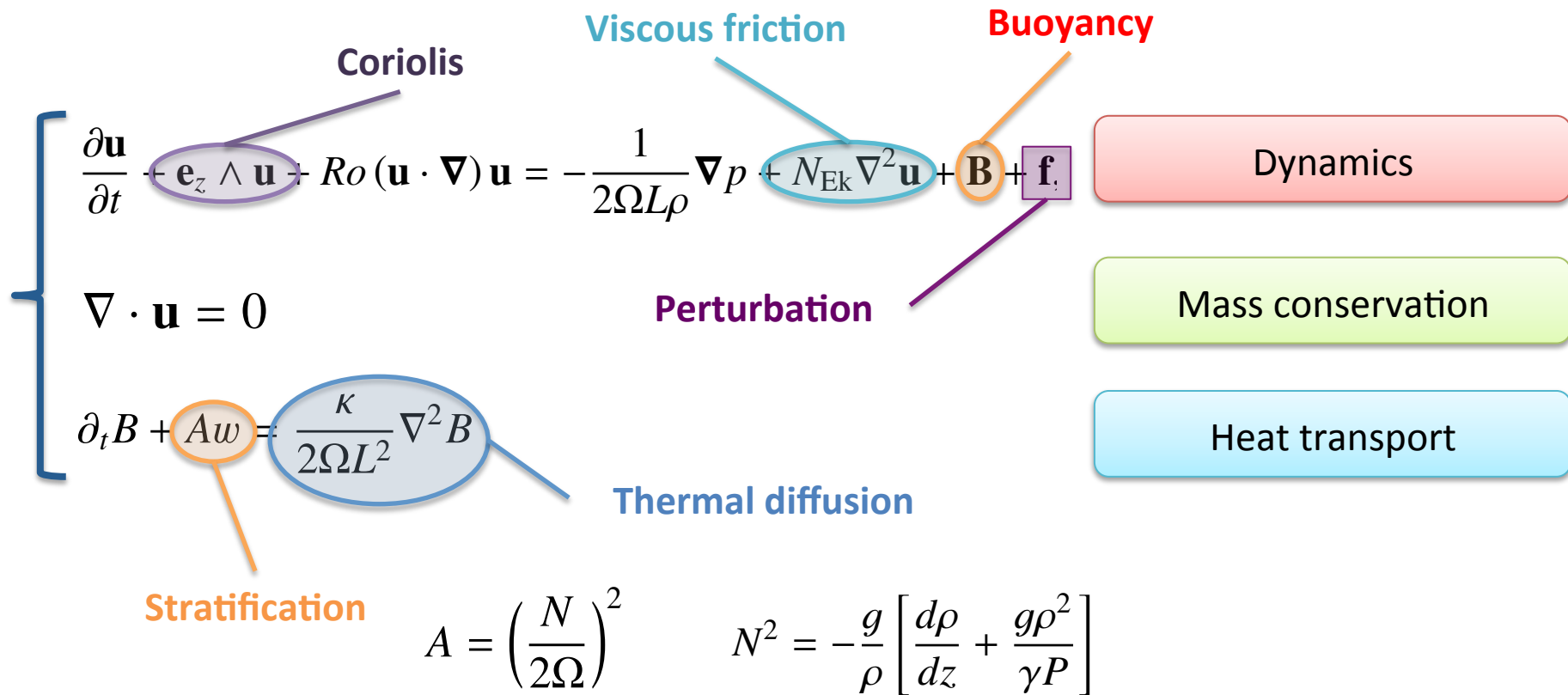


Fig. Ogilvie & Lin (2004)



III. How does tidal dissipation depend on the internal physics of the bodies ?

- Equations



III. How does tidal dissipation depend on the internal physics of the bodies ?

- Control parameters

θ

Colatitude

Position

$$A = \left(\frac{N}{2\Omega} \right)^2$$

Frequency ratio

Nature of the waves

$$E = \frac{2\pi^2\nu}{\Omega L^2}$$

Ekman

Influence of viscous friction compared to Coriolis effects

$$K = \frac{2\pi^2\kappa}{\Omega L^2}$$

Thermal diffusivity

Influence of thermal diffusion compared to Coriolis effects

III. How does tidal dissipation depend on the internal physics of the bodies ?

- Periodic velocity field $u = \sum u_{mn} e^{i2\pi(mx+nz)}$

$$\left\{ \begin{array}{l}
 u_{mn} = n \frac{i\tilde{\omega} (nf_{mn} - mh_{mn}) - n \cos \theta g_{mn}}{(m^2 + n^2) \tilde{\omega}^2 - n^2 \cos^2 \theta - Am^2 \frac{\tilde{\omega}}{\hat{\omega}}} \\
 v_{mn} = \frac{n \cos \theta (nf_{mn} - mh_{mn}) + i \left[(m^2 + n^2) \tilde{\omega} - \frac{Am^2}{\hat{\omega}} \right] g_{mn}}{(m^2 + n^2) \tilde{\omega}^2 - n^2 \cos^2 \theta - Am^2 \frac{\tilde{\omega}}{\hat{\omega}}} \\
 w_{mn} = -m \frac{i\tilde{\omega} (nf_{mn} - mh_{mn}) - n \cos \theta g_{mn}}{(m^2 + n^2) \tilde{\omega}^2 - n^2 \cos^2 \theta - Am^2 \frac{\tilde{\omega}}{\hat{\omega}}}
 \end{array} \right.$$

Influence of the perturbation
Inertial part

Viscous diffusivity

$$\left\{ \begin{array}{l}
 \tilde{\omega} = \omega + iE(m^2 + n^2) \\
 \hat{\omega} = \omega + iK(m^2 + n^2)
 \end{array} \right.$$

Thermal diffusivity

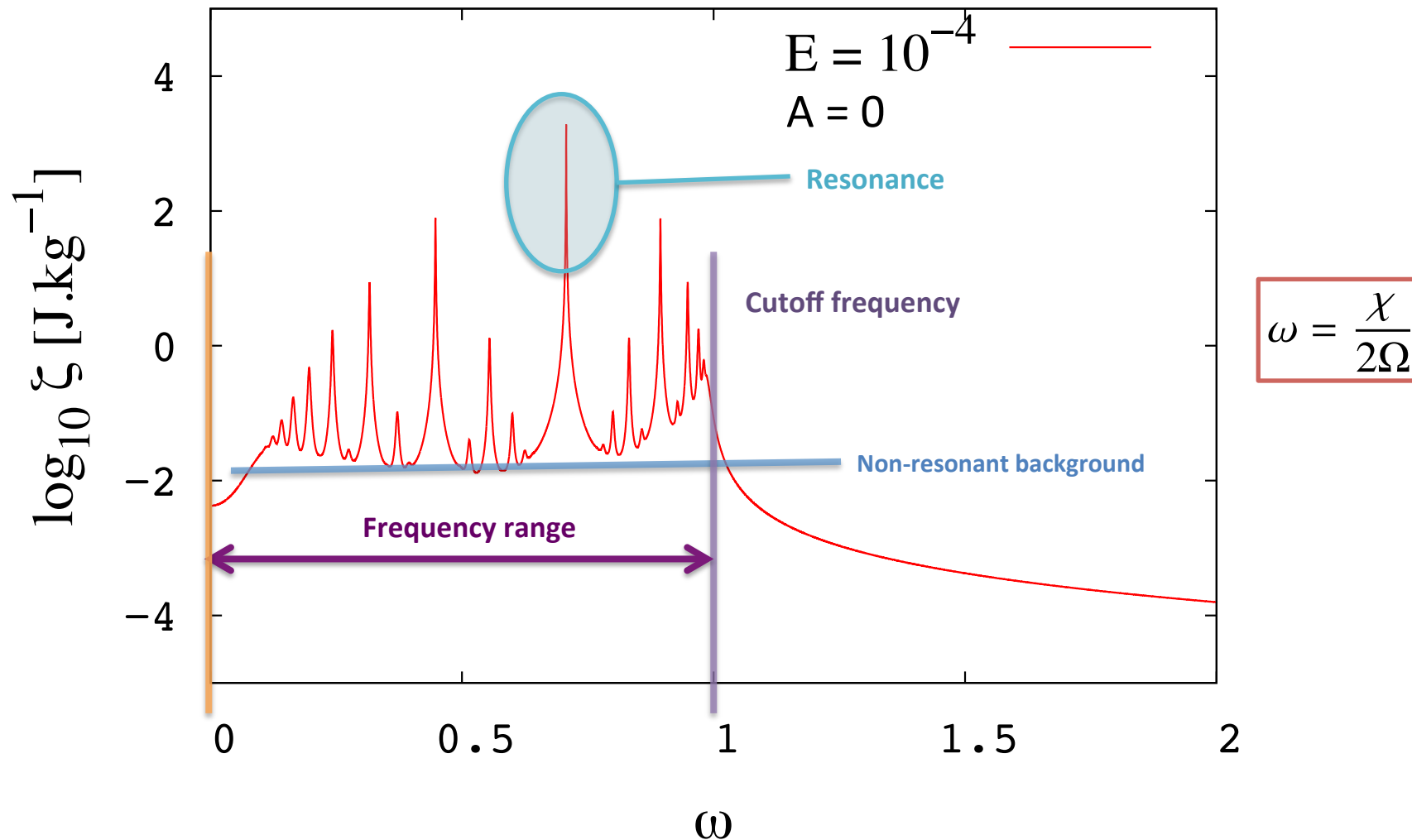
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Energy dissipated by viscous friction

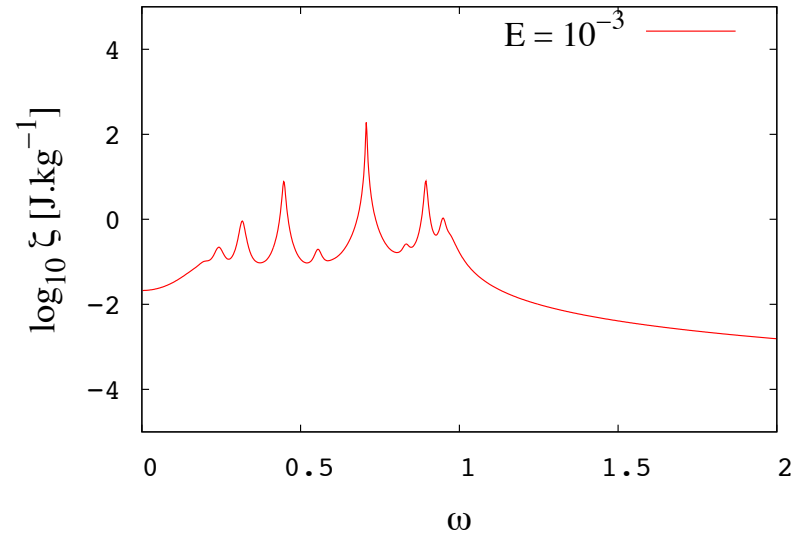
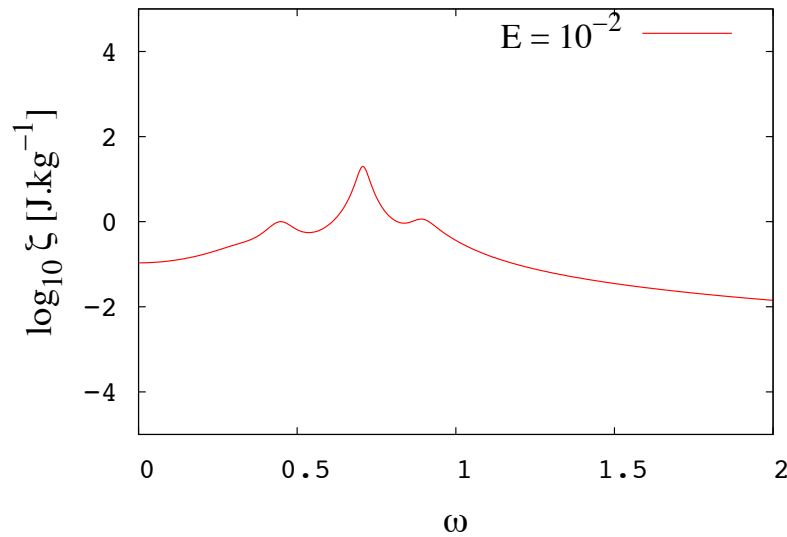
$$D = \int_0^1 \int_0^1 \langle -\mathbf{u} \cdot \nu \nabla^2 \mathbf{u} \rangle dx dz$$

$$\zeta = \frac{2\pi}{\Omega} D = 2\pi E \sum_{(m,n) \in \mathbb{Z}^{*2}} (m^2 + n^2) (|u_{mn}^2| + |v_{mn}^2| + |w_{mn}^2|)$$

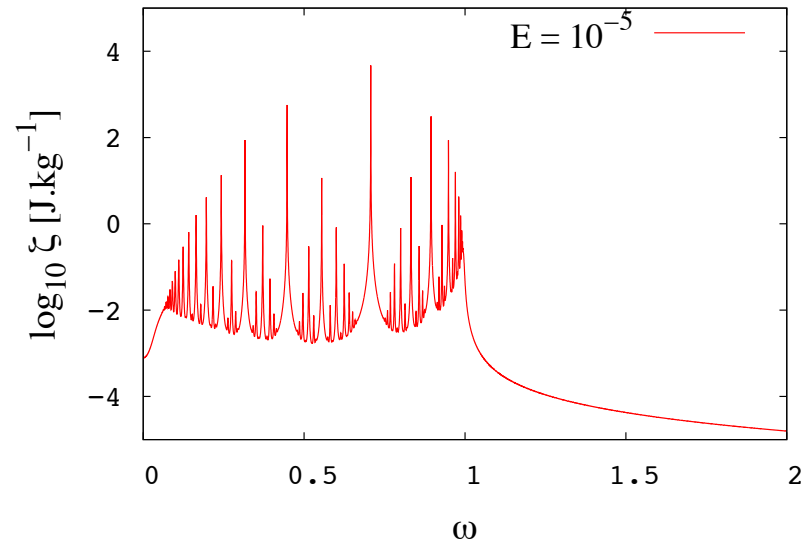
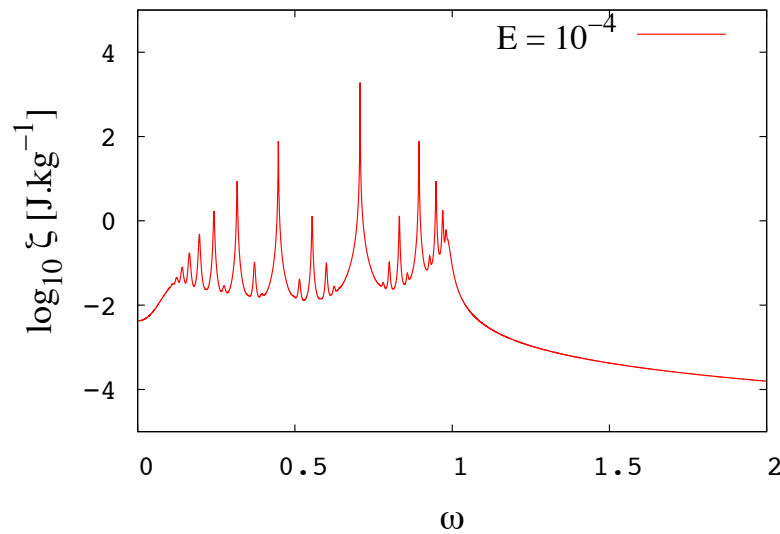
III. How does tidal dissipation depend on the internal physics of the bodies ?



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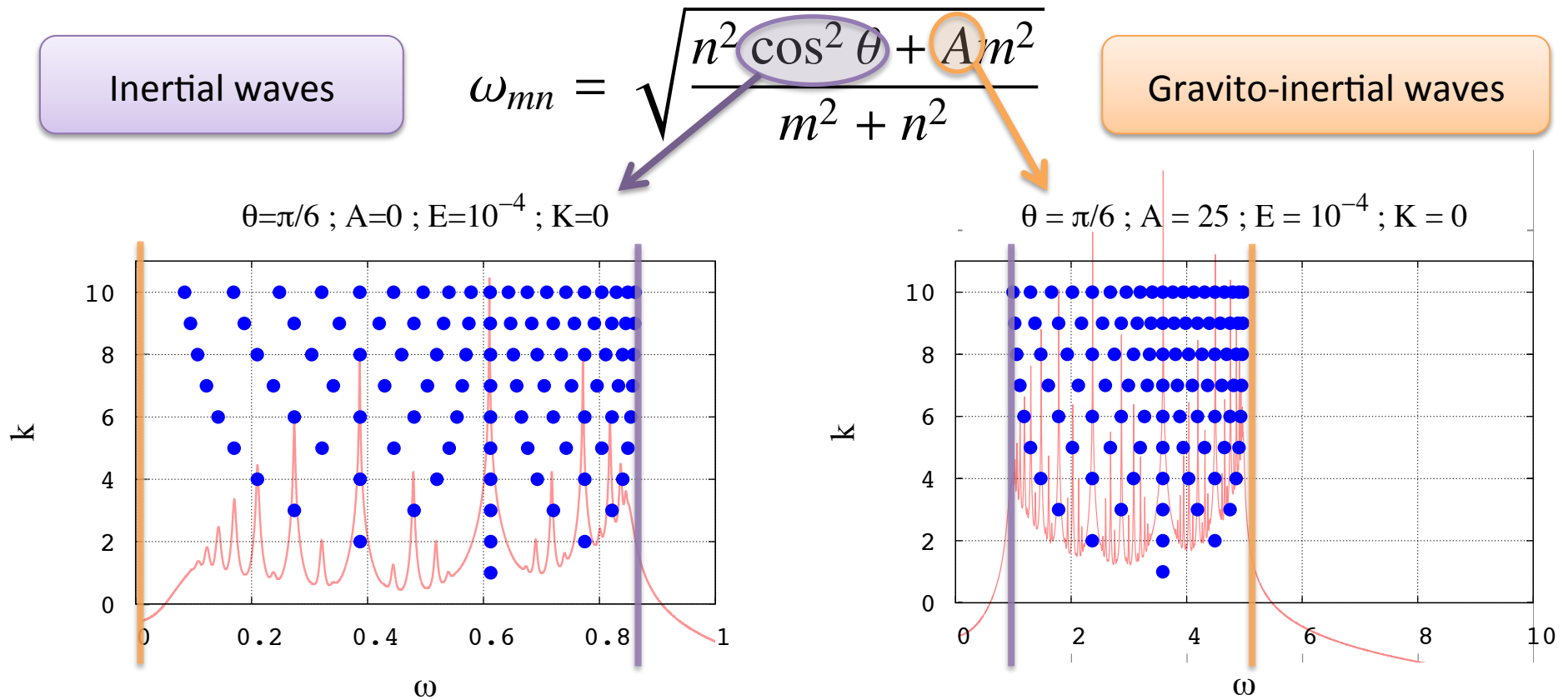


$\theta = 0$
 $A = 0$
 $K = 0$



III. How does tidal dissipation depend on the internal physics of the bodies ?

- Positions of the resonances

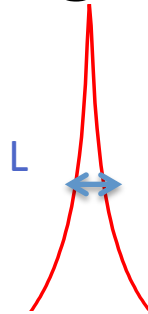


$$k = \max \{|m|, |n|\}$$

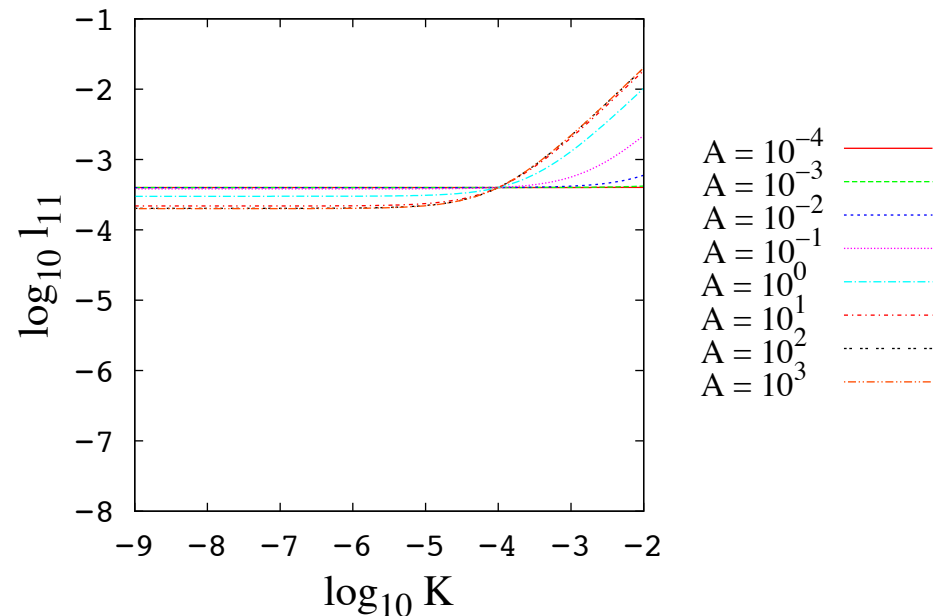
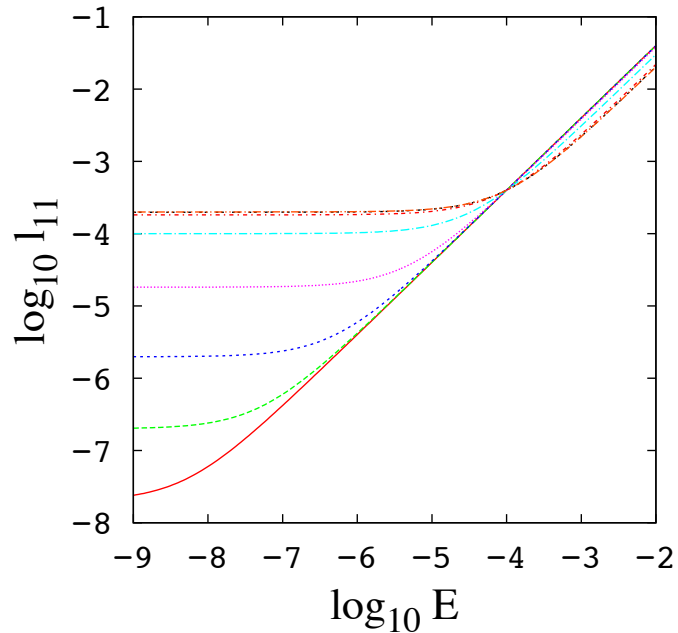
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Width at mid-height

$$l_{mn} = (m^2 + n^2) \frac{Am^2K + (2n^2 \cos^2 \theta + Am^2)E}{n^2 \cos^2 \theta + Am^2}$$



DOMAIN	$A \ll A_{mn}$	$A \gg A_{mn}$
$Pr \gg Pr_{mn}$	$2E(m^2 + n^2)$	$E(m^2 + n^2)$
$Pr \ll Pr_{mn}$	$AK \frac{m^2(m^2 + n^2)}{n^2 \cos^2 \theta}$	$K(m^2 + n^2)$

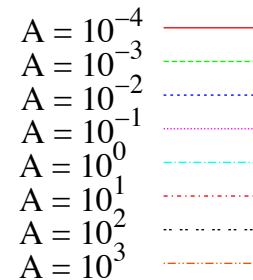
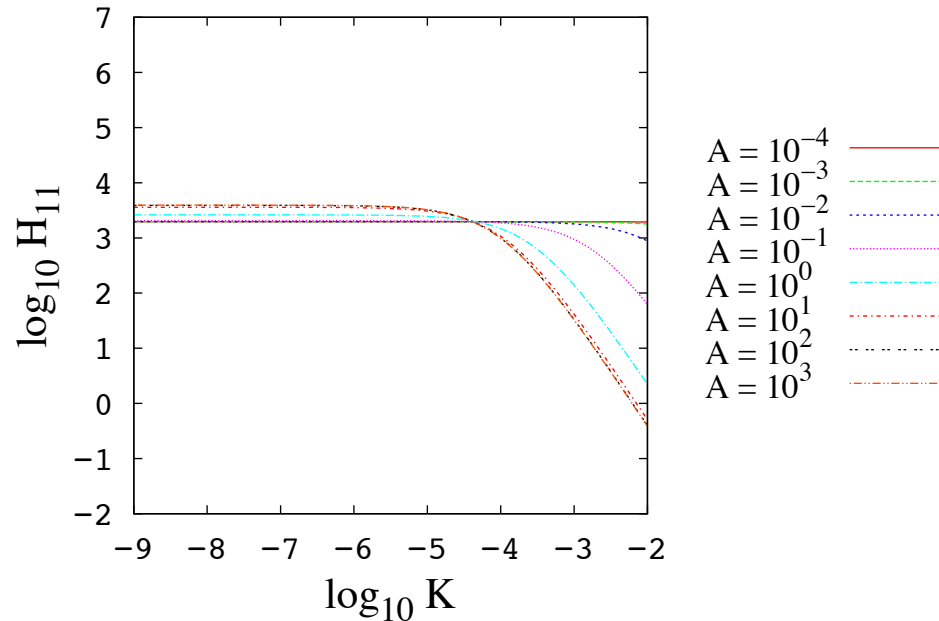
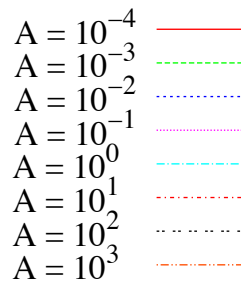
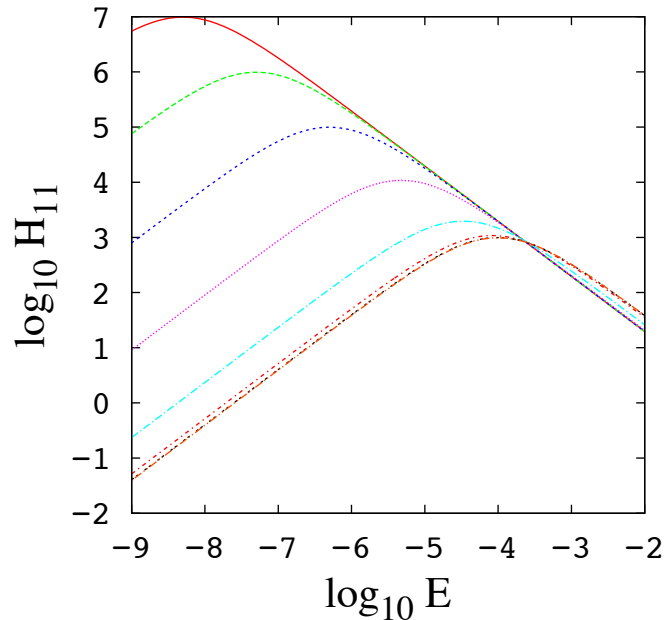


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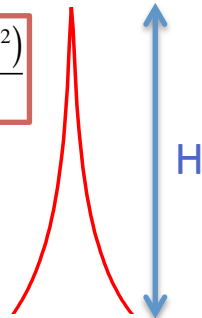
- Height

$$H_{mn} = \frac{8\pi F^2 E}{m^2 n^2 (m^2 + n^2)^2} \frac{(2n^2 \cos^2 \theta + Am^2)(n^2 \cos^2 \theta + Am^2)}{[Am^2 K + (2n^2 \cos^2 \theta + Am^2) E]^2}$$

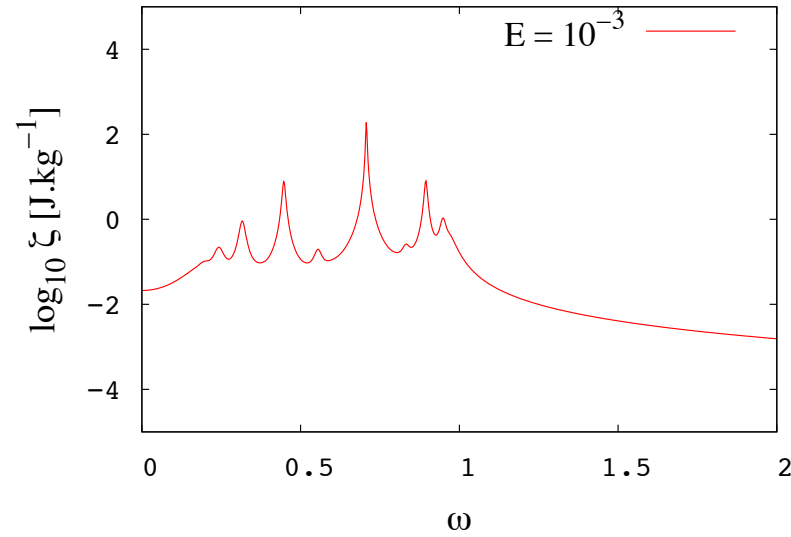
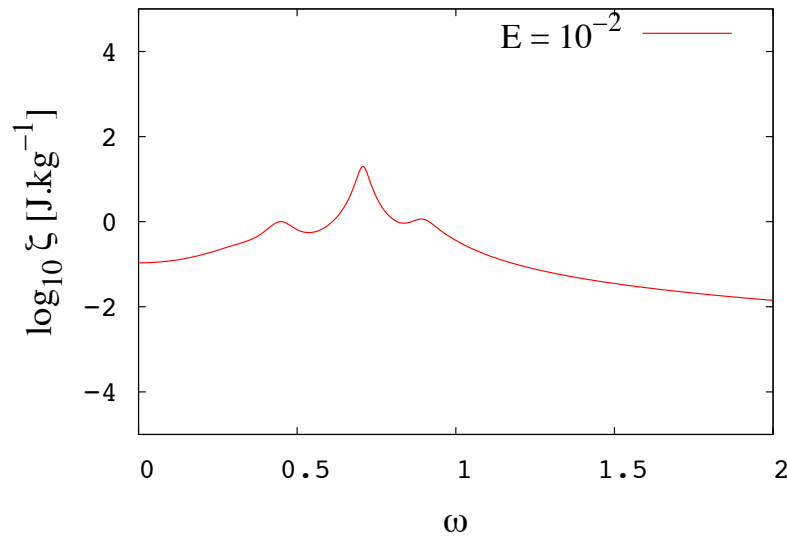
$$f_{mn} = i \frac{F}{|m| n^2}, \quad g_{mn} = 0, \quad h_{mn} = 0$$



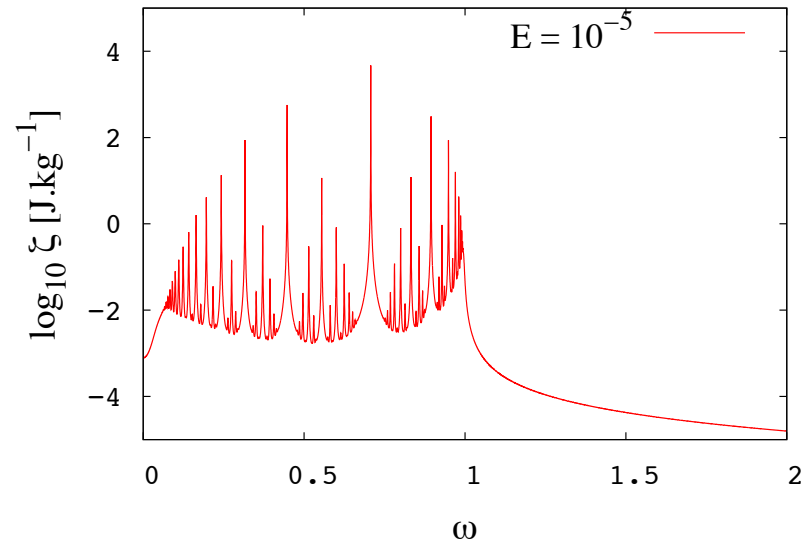
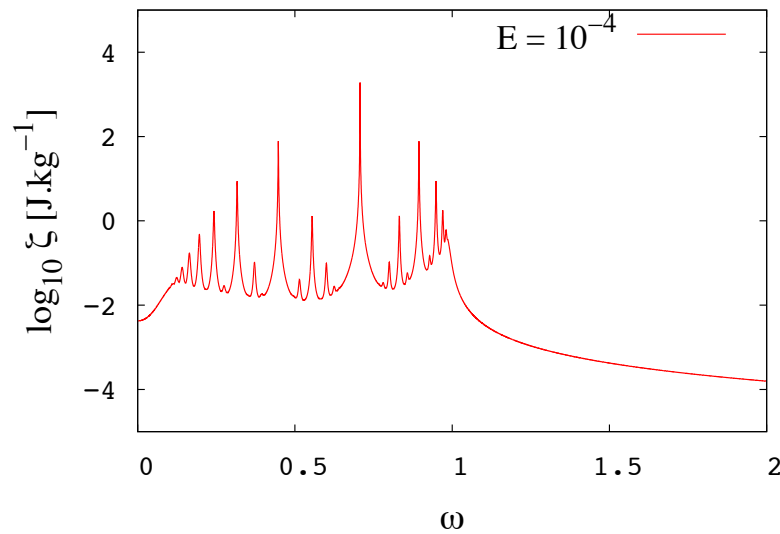
DOMAIN	$A \ll A_{mn}$	$A \gg A_{mn}$
$Pr \gg Pr_{mn}$	$\frac{4\pi F^2}{m^2 n^2 (m^2 + n^2)^2 E}$	$\frac{8\pi F^2}{m^2 n^2 (m^2 + n^2)^2 E}$
$Pr \ll Pr_{mn}$	$\frac{16\pi F^2 n^2}{m^6 (m^2 + n^2)^2} \frac{E \cos^4 \theta}{A^2 K^2}$	$\frac{8\pi F^2}{m^2 n^2 (m^2 + n^2)^2} \frac{E}{K^2}$



III. How does tidal dissipation depend on the internal physics of the bodies ?



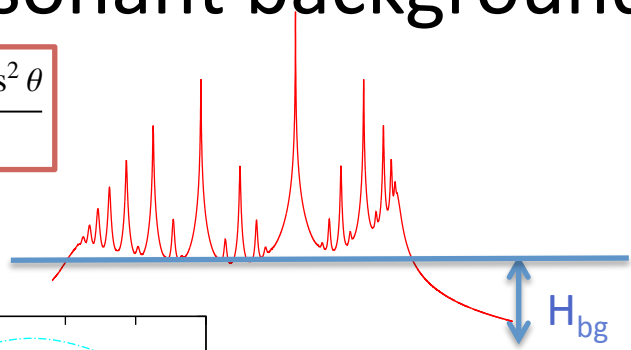
$\theta = 0$
 $A = 0$
 $K = 0$



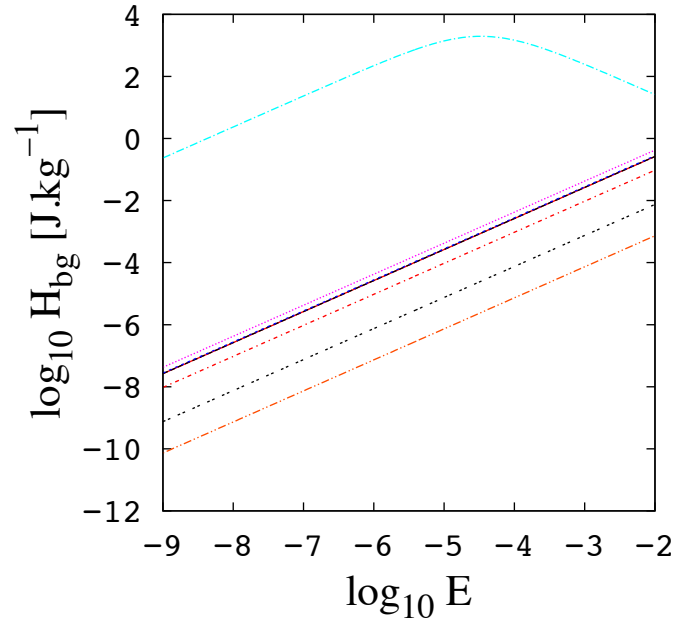
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Non-resonant background

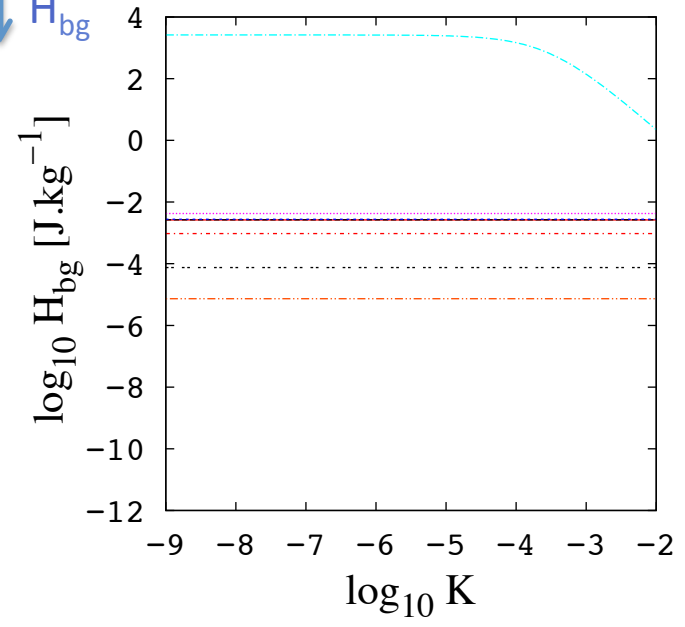
$$H_{\text{bg}} = 4\pi F^2 E \frac{C_{\text{grav}}^\infty A + C_{\text{in}}^\infty \cos^2 \theta}{(A + \cos^2 \theta)^2}$$



$A \ll \cos^2 \theta$	$A \gg \cos^2 \theta$
$4\pi C_{\text{in}}^\infty F^2 \frac{E}{\cos^2 \theta}$	$4\pi C_{\text{grav}}^\infty F^2 \frac{E}{A}$

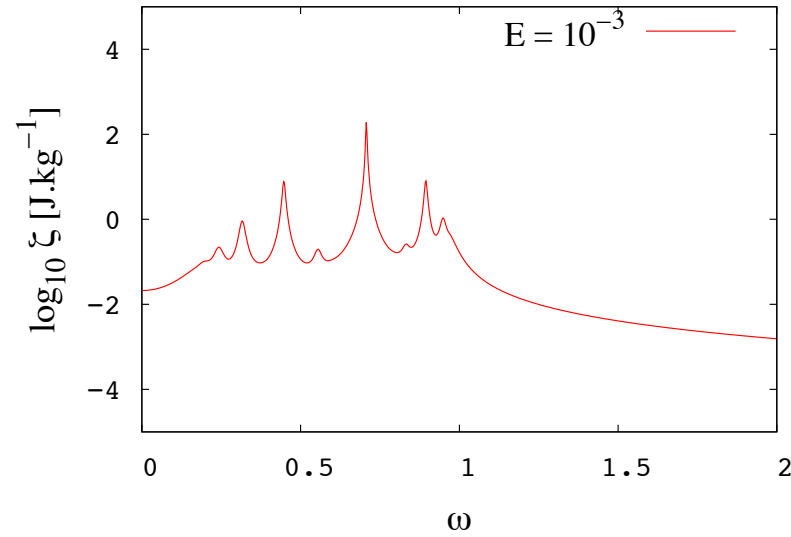
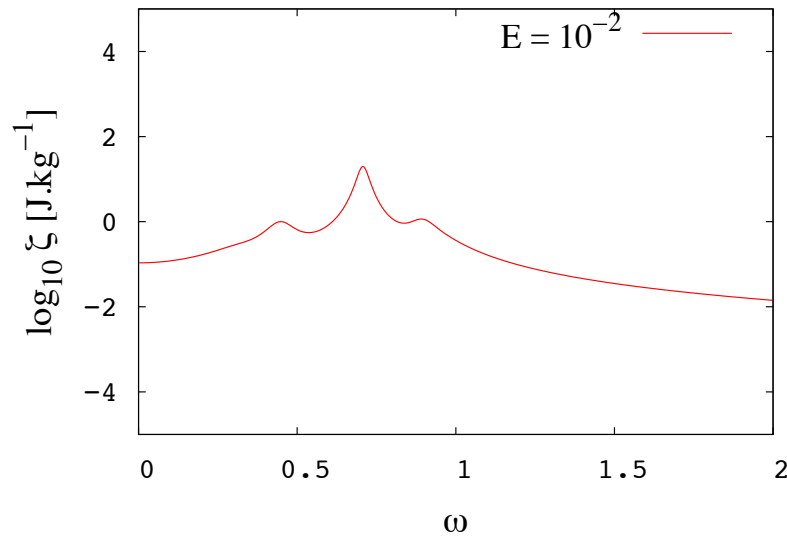


- $A = 10^{-4}$ ———
- $A = 10^{-3}$ - - - -
- $A = 10^{-2}$ ·····
- $A = 10^{-1}$ -·-·-
- $A = 10^0$ ———
- $A = 10^1$ -·-·-
- $A = 10^2$ - - - -
- $A = 10^3$ ———

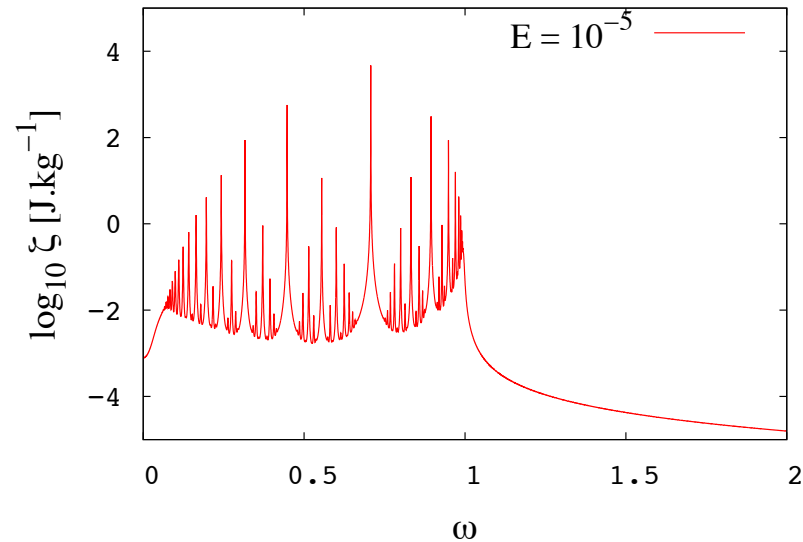
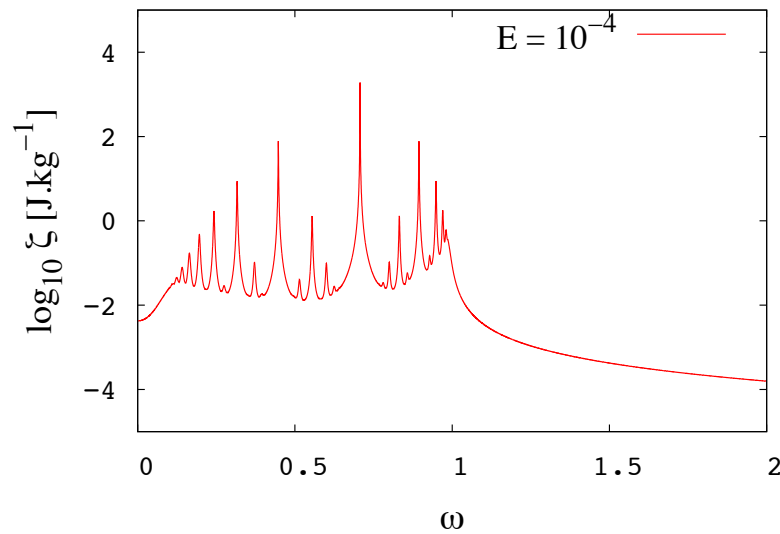


- $A = 10^{-4}$ ———
- $A = 10^{-3}$ - - - -
- $A = 10^{-2}$ ·····
- $A = 10^{-1}$ -·-·-
- $A = 10^0$ ———
- $A = 10^1$ -·-·-
- $A = 10^2$ - - - -
- $A = 10^3$ ———

III. How does tidal dissipation depend on the internal physics of the bodies ?



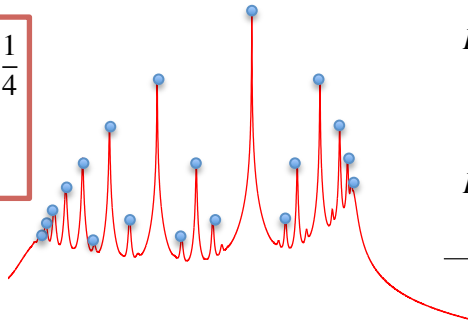
$\theta = 0$
 $A = 0$
 $K = 0$



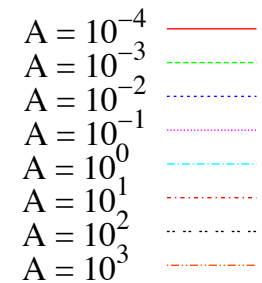
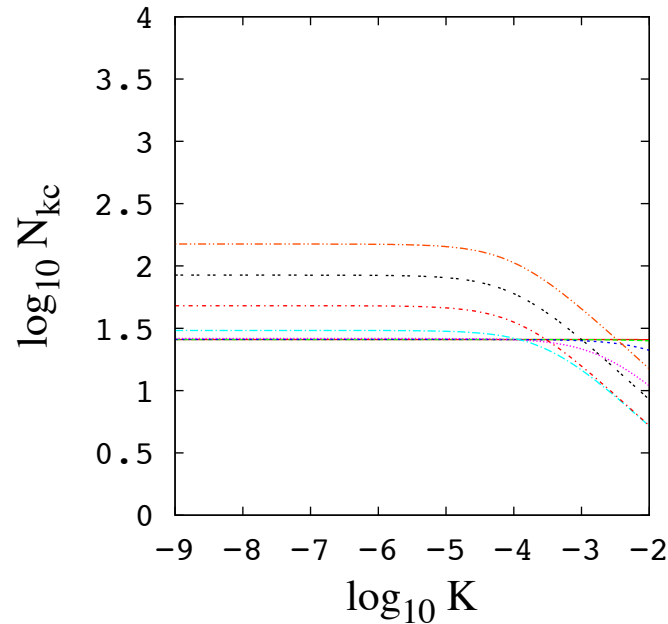
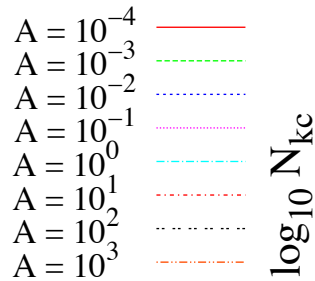
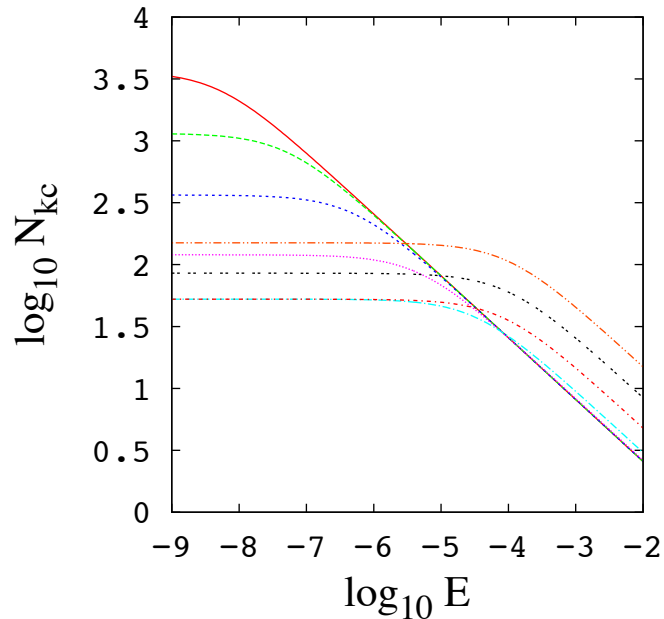
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Number of resonances

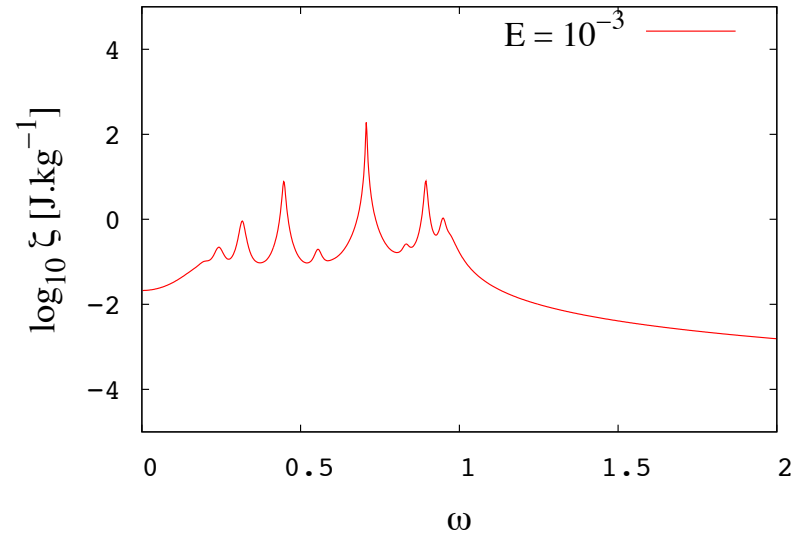
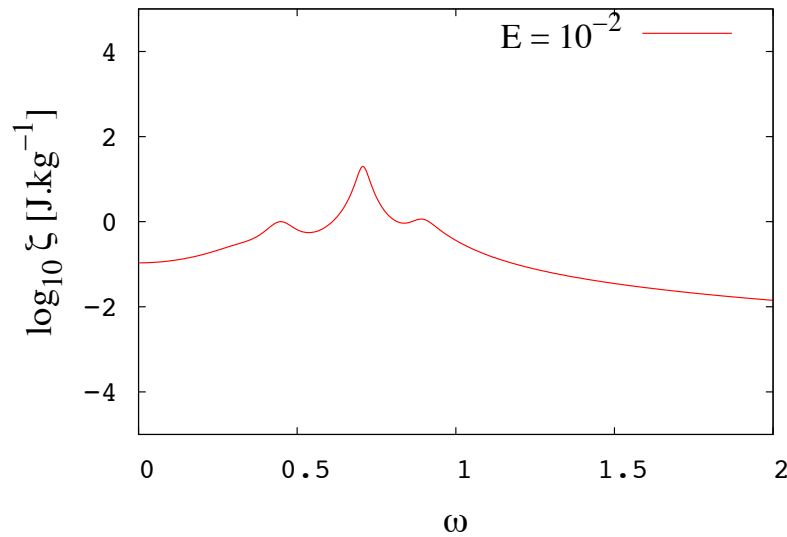
$$N_{kc} \sim \left\{ \frac{1}{2} \frac{(2 \cos^2 \theta + A)(A + \cos^2 \theta)^3}{[AK + (2 \cos^2 \theta + A)E]^2 [C_{in}^\infty \cos^2 \theta + C_{grav}^\infty A]} \right\}^{\frac{1}{4}}$$



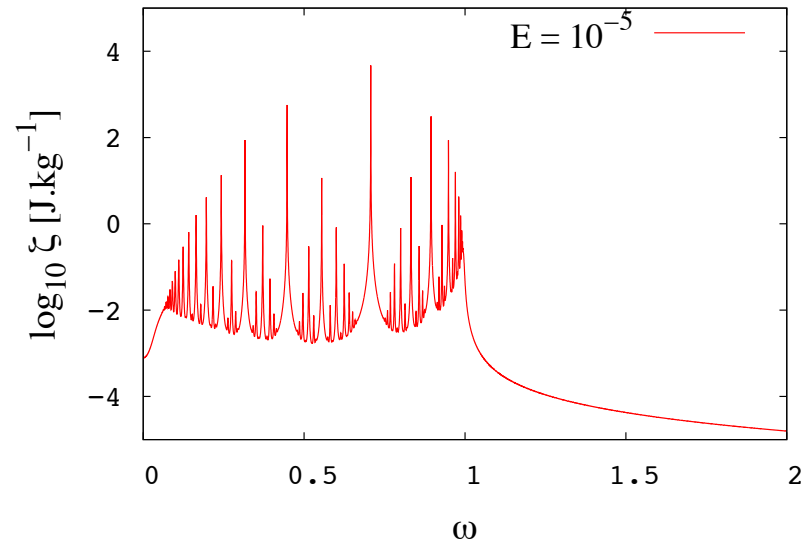
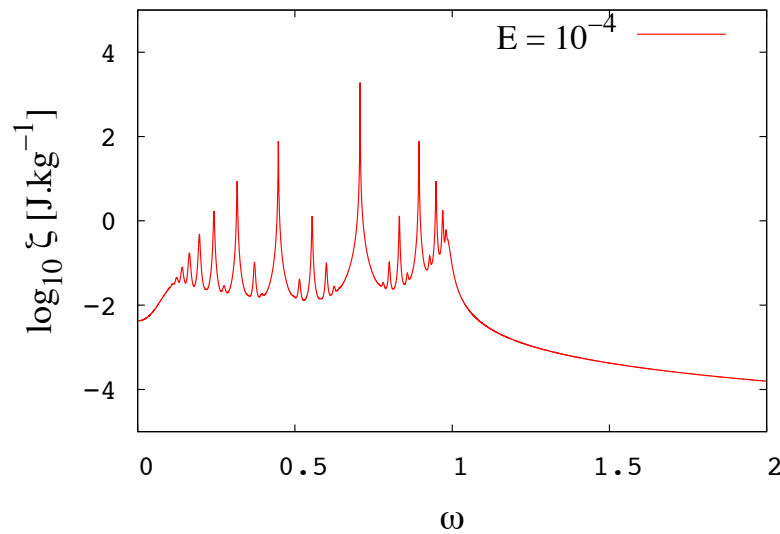
DOMAIN	$A \ll A_{11}$	$A \gg A_{11}$
$Pr \gg Pr_{11}$	$N_{kc} \sim \left(\frac{\cos^2 \theta}{4C_{in}^\infty E^2} \right)^{\frac{1}{4}}$	$N_{kc} \sim \left(\frac{A}{2C_{grav}^\infty E^2} \right)^{\frac{1}{4}}$
$Pr \ll Pr_{11}$	$N_{kc} \sim \left(\frac{\cos^6 \theta}{C_{in}^\infty A^2 K^2} \right)^{\frac{1}{4}}$	$N_{kc} \sim \left(\frac{A}{2C_{grav}^\infty K^2} \right)^{\frac{1}{4}}$



III. How does tidal dissipation depend on the internal physics of the bodies ?



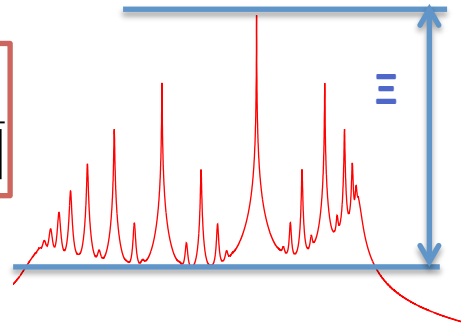
$\theta = 0$
 $A = 0$
 $K = 0$



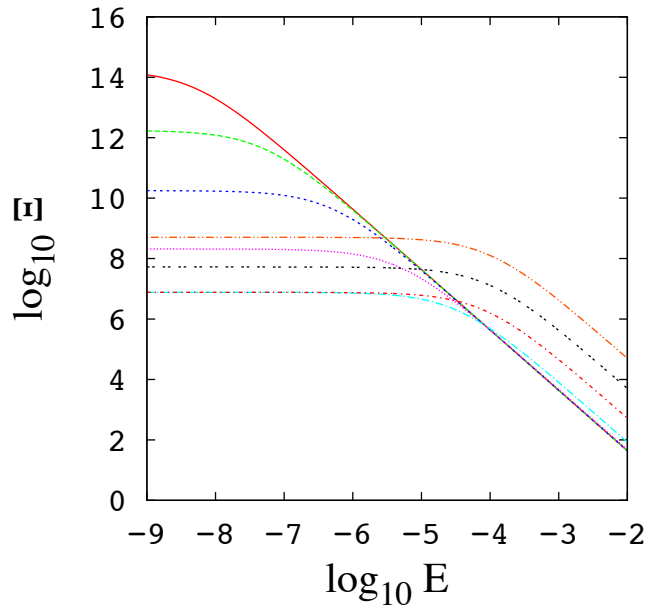
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Sharpness ratio

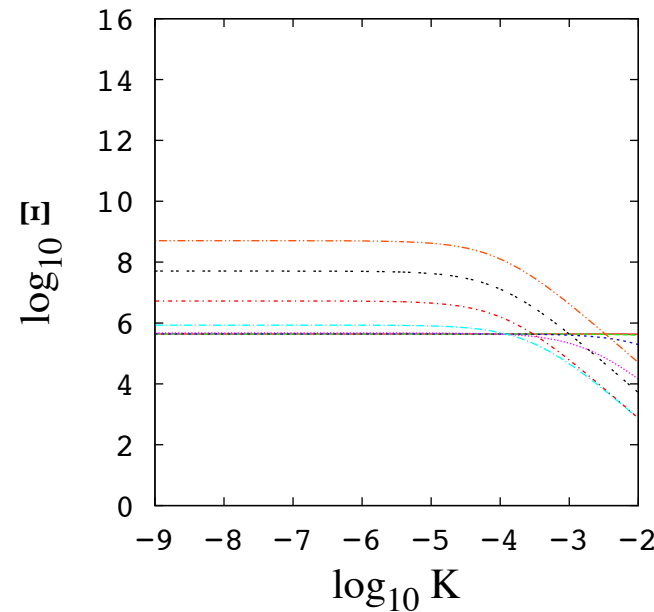
$$\Xi = \frac{1}{2} \frac{(2 \cos^2 \theta + A)(A + \cos^2 \theta)^3}{[AK + (2 \cos^2 \theta + A)E]^2 [C_{\text{in}}^\infty \cos^2 \theta + C_{\text{grav}}^\infty A]}$$



DOMAIN	$A \ll A_{11}$	$A \gg A_{11}$
$Pr \gg Pr_{11}$	$\frac{\cos^2 \theta}{4C_{\text{in}}^\infty E^2}$	$\frac{A}{2C_{\text{grav}}^\infty E^2}$
$Pr \ll Pr_{11}$	$\frac{\cos^6 \theta}{C_{\text{in}}^\infty A^2 K^2}$	$\frac{A}{2C_{\text{grav}}^\infty K^2}$

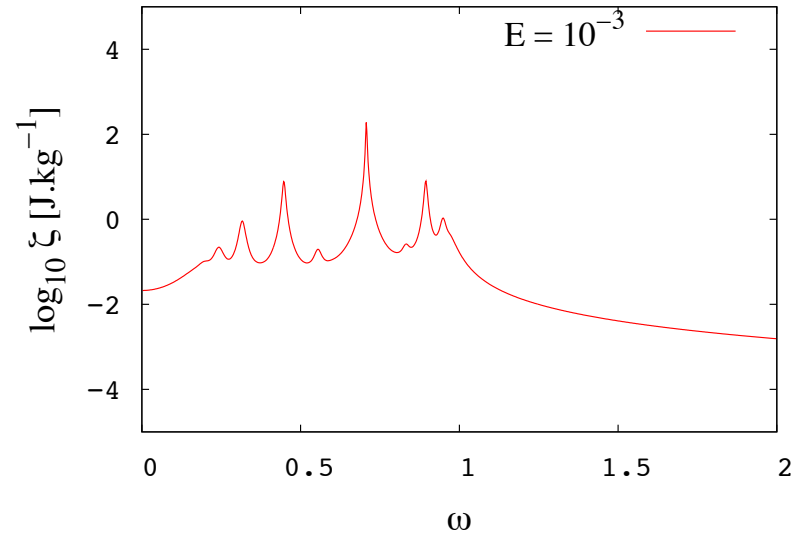
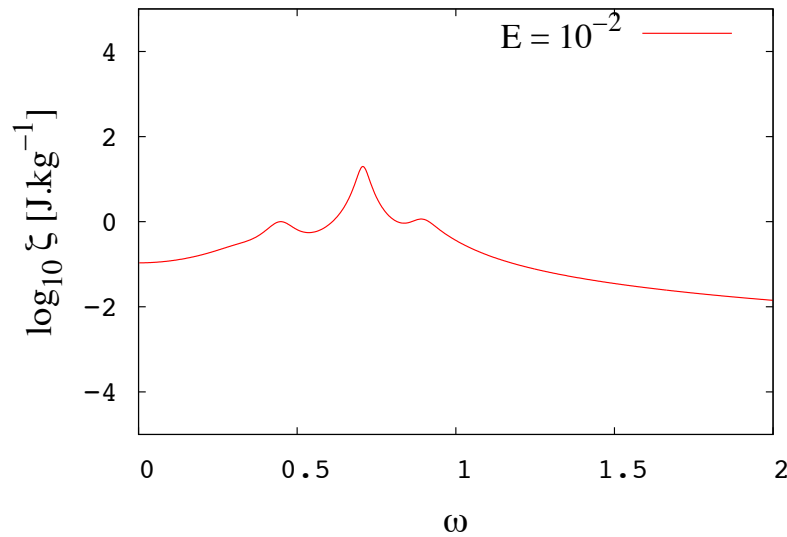


$A = 10^{-4}$ ———
 $A = 10^{-3}$ - - -
 $A = 10^{-2}$ ····
 $A = 10^{-1}$ ····
 $A = 10^0$ - - -
 $A = 10^1$ ····
 $A = 10^2$ ····
 $A = 10^3$ - - -

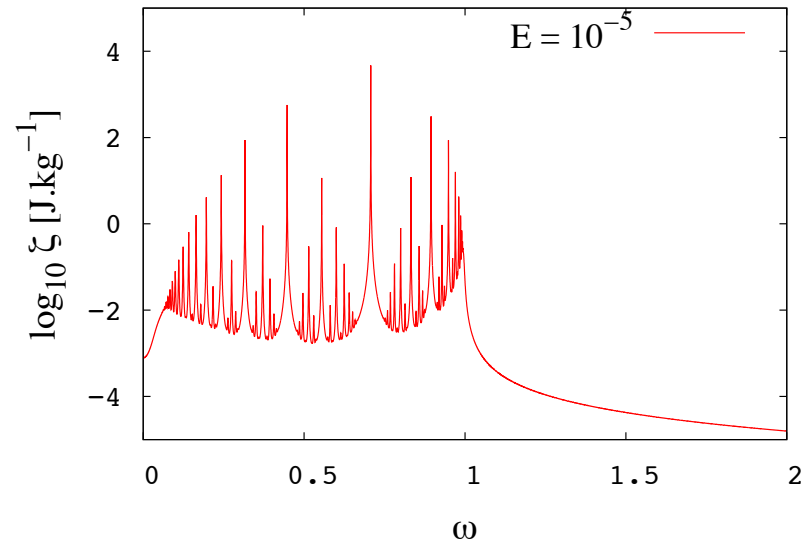
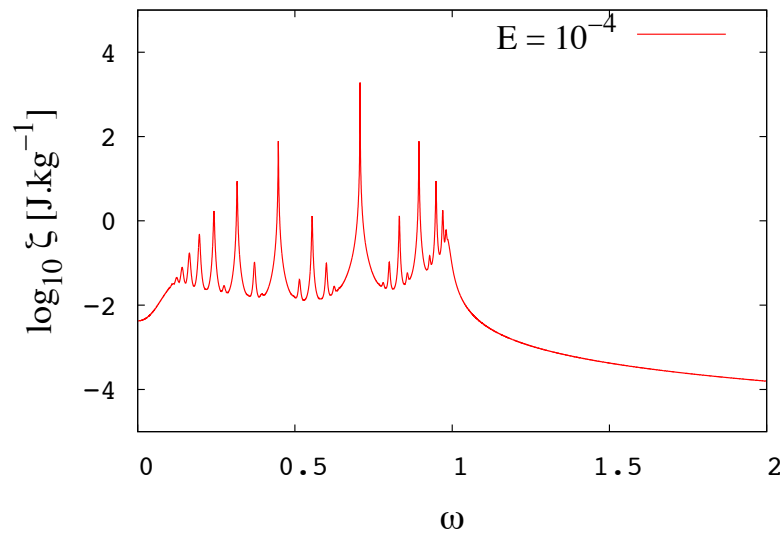


$A = 10^{-4}$ ———
 $A = 10^{-3}$ - - -
 $A = 10^{-2}$ ····
 $A = 10^{-1}$ ····
 $A = 10^0$ - - -
 $A = 10^1$ ····
 $A = 10^2$ ····
 $A = 10^3$ - - -

III. How does tidal dissipation depend on the internal physics of the bodies ?



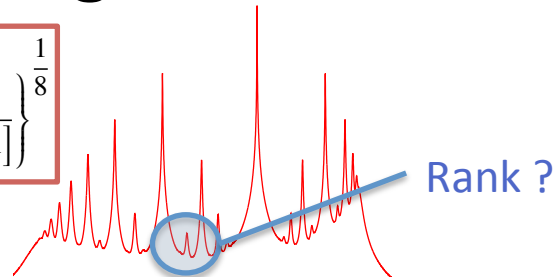
$\theta = 0$
 $A = 0$
 $K = 0$



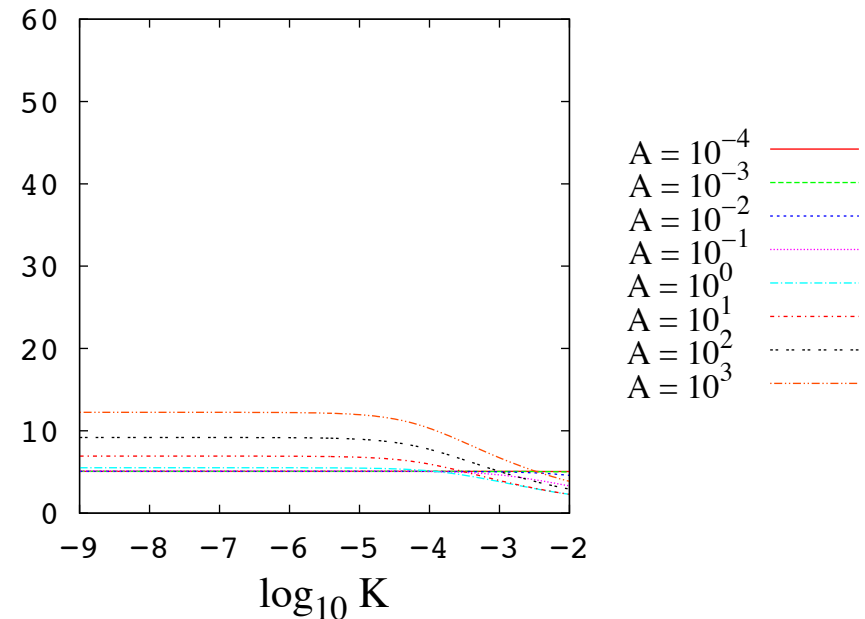
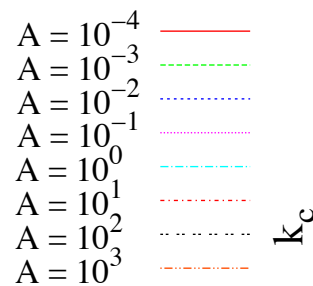
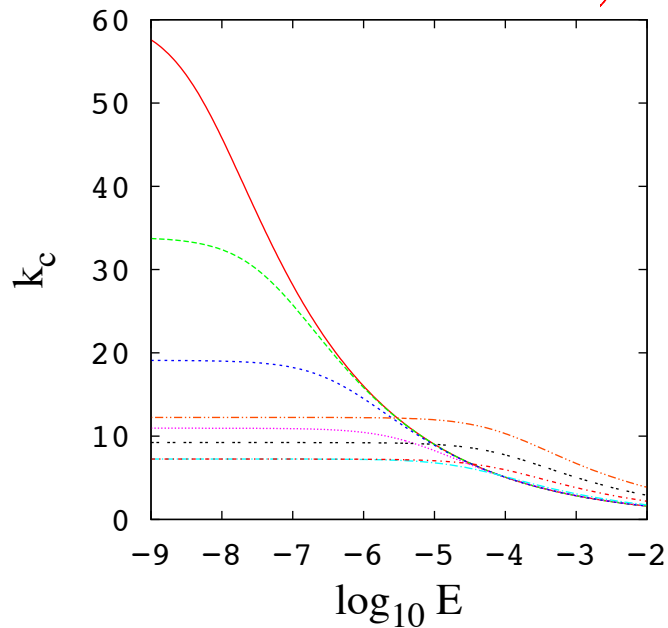
III. How does tidal dissipation depend on the internal physics of the bodies ?

- Rank of the highest harmonics

$$k_c \sim \left\{ \frac{1}{2} \frac{(2 \cos^2 \theta + A)(A + \cos^2 \theta)^3}{[AK + (2 \cos^2 \theta + A)E]^2 [C_{in}^\infty \cos^2 \theta + C_{grav}^\infty A]} \right\}^{\frac{1}{8}}$$



DOMAIN	$A \ll A_{11}$	$A \gg A_{11}$
$Pr \gg Pr_{11}$	$k_c \sim \left(\frac{\cos^2 \theta}{4C_{in}^\infty E^2} \right)^{\frac{1}{8}}$	$k_c \sim \left(\frac{A}{2C_{grav}^\infty E^2} \right)^{\frac{1}{8}}$
$Pr \ll Pr_{11}$	$k_c \sim \left(\frac{\cos^6 \theta}{C_{in}^\infty A^2 K^2} \right)^{\frac{1}{8}}$	$k_c \sim \left(\frac{A}{2C_{grav}^\infty K^2} \right)^{\frac{1}{8}}$



Synopsis

- I. Tidal dissipation ?
- II. Tidal gravito-inertial waves ?
- III. Impact of tidal dissipation on the spin/orbital evolution ?
- IV. How does tidal dissipation depend on the internal physics of the bodies ?

